

Optimizing the location of the ^{13}C inversion pulse in a modified INEPT sequence: Illustration with $^{13}\text{C}_3$ glutamate

Atiyah Yahya^{1,2} and Peter S. Allen³

¹Department of Medical Physics, Cross Cancer Institute, Edmonton, Alberta, Canada, ²Department of Oncology, University of Alberta, Edmonton, Alberta, Canada,

³Department of Biomedical Engineering, University of Alberta, Edmonton, Alberta, Canada

Target Audience: Researchers working in the field of ^{13}C NMR (Nuclear Magnetic Resonance) spectroscopy.

Purpose: To optimize the signal to noise ratio (SNR) of ^{13}C glutamate (Glu) signal obtained with the localized INEPT (Insensitive Nuclei Enhanced by Polarization Transfer) sequence, also referred to as LINEPT¹. The INEPT sequence shown in Fig. 1(a) allows for the direct detection of enhanced ^{13}C signal from ^1H -J-coupled ^{13}C nuclei. Spatial localization was incorporated into the sequence by replacing the proton hard pulses by slice selective ones and separating the ^{13}C inversion pulse from the proton refocusing pulse in time to allow the echo time (TE) to be longer than $1/(2J_{\text{CH}})$ ¹, which is only ≈ 3.7 ms for the C₃ and C₄ carbons of Glu (J_{CH} is the heteronuclear coupling constant). The repositioning of the ^{13}C inversion pulse was based on heteronuclear product operator calculations which indicate that whether the ^{13}C 180° pulse is placed 1/(4 J_{CH}) after the first proton pulse or 1/(4 J_{CH}) before the third proton pulse, the required heteronuclear coupling evolution for polarization transfer will occur (Fig. 1(b)). The objective of this work is to demonstrate, using $^{13}\text{C}_3$ -Glu at 3 T as an example, that the concurrent strong proton homonuclear coupling evolutions taking place causes the ^{13}C signal yield of a refocused modified INEPT sequence (Fig. 1(c)) to be dependent on whether the ^{13}C inversion pulse is placed at the beginning or at the end of the sequence. Our results show that the signal, which also varies with TE, is significantly higher if the pulse is placed 1/(4 J_{CH}) after the proton excitation pulse.

Methods: Numerical density matrix calculations were performed in MATLAB by representing $^{13}\text{C}_3$ -Glu as an AMNPQX spin system. The MN protons are strongly-coupled to each other and to the PQ protons and weakly-coupled to the $^{13}\text{C}_3$ nuclei ($J_{\text{CH}} = 135$ Hz). The sequence of Figure 1(c) was implemented on a whole-body 3 T scanner; hard pulses were employed for simplicity and a refocused version of INEPT was employed to allow the formation of in-phase ^{13}C signal (to enable proton decoupling). Two versions of the sequence were employed, one with the ^{13}C inversion pulse applied 1/(4 J_{CH}) ≈ 1.85 ms after the proton excitation pulse, and the other with the pulse placed 1.85 ms prior to the third proton pulse. Experiments were conducted on a 2 cm diameter phantom containing 35 mM of 99% enriched $^{13}\text{C}_3$ -Glu with a home-built 7 cm diameter ^1H RF (radiofrequency) coil and a 3.5 cm diameter ^{13}C RF coil. TE values were varied between 10 and 30 ms in steps of 5 ms and a TR (repetition time) of 3 seconds was employed. WALTZ-16 proton decoupling was applied during acquisition and a 16-step phase cycling scheme was used.

Results: Figure 2 shows the numerically calculated signal intensity of the decoupled $^{13}\text{C}_3$ – Glu peak as a function of TE in response to the two versions of the sequence displayed in Fig. 1(c). Intensities are normalized to the higher TE = 10 ms value. It can be seen that the signal decreases with increasing TE; furthermore, the signal is always higher when the ^{13}C inversion pulse is placed 1/(4 J_{CH}) after the first proton pulse. For example, when TE = 30 ms the signal decreases by more than 50% when the ^{13}C 180° is applied towards the end of the sequence. Experimental data points acquired with five TE values agree with the numerical calculations. A similar enhancement was found numerically for $^{13}\text{C}_4$ – Glu.

Discussion: The $^{13}\text{C}_3$ -Glu signal is governed by the amount of heteronuclear proton antiphase coherences $2M_XX_Z$ and $2N_XX_Z$ present immediately prior to the two 90° simultaneous pulses responsible for polarization transfer. It was calculated that the expectation value of the mentioned coherences follow the same trends shown in Fig. 2. With increasing TE, the antiphase coherences decrease due to a loss of signal to other Glu protons as a result of strong homonuclear coupling². The non-intuitive decline of the coherences when the ^{13}C inversion pulse is placed towards the end of the sequence was also found to be due to the strong coupling between the MN and PQ protons. The amount of the heteronuclear antiphase M and N coherences are formed not only from in-phase M_Y and N_Y magnetization (formed by the ^1H excitation pulse) but also from in-phase P_Y and Q_Y magnetization. We found that the contribution from the latter declines when the ^{13}C inversion pulse is placed 1/(4 J_{CH}) before the third proton pulse causing the final observed ^{13}C signal to be lower.

Conclusion: To obtain higher $^{13}\text{C}_3$ or $^{13}\text{C}_4$ -glutamate signal from a modified localized INEPT (LINEPT) sequence, the ^{13}C inversion pulse should be placed 1/(4 J_{CH}) after the first proton pulse rather than 1/(4 J_{CH}) before the third proton pulse.

References:

1. Watanabe H, Ishihara Y, Okamoto K et al. *In Vivo* 3D Localized ^{13}C Spectroscopy Using Modified INEPT and DEPT. *J. Magn. Reson.* 1998; 134:214-222.
2. Kay LE, McClung RED, A Product Operator Description of AB and ABX Spin Systems. *J. Magn. Reson.* 1988; 77: 258-273

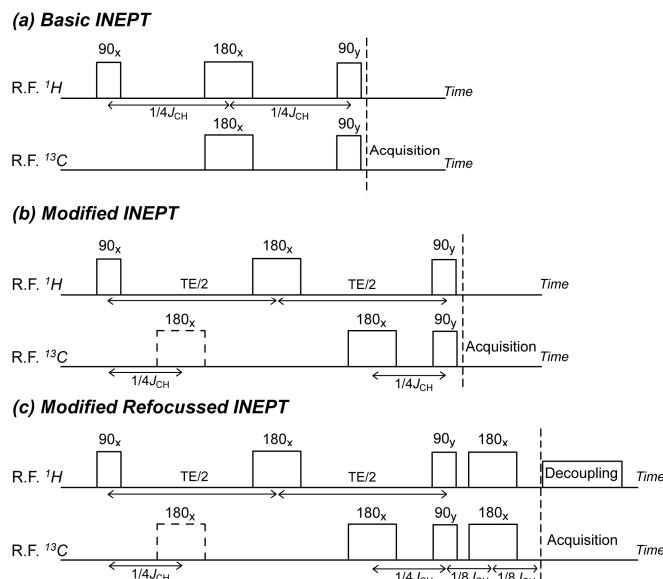


Figure 1: The modified INEPT sequence¹ separates the ^{13}C inversion pulse from the proton refocusing pulse in time placing it either 1/(4 J_{CH}) after the first proton pulse or 1/(4 J_{CH}) before the third proton pulse.

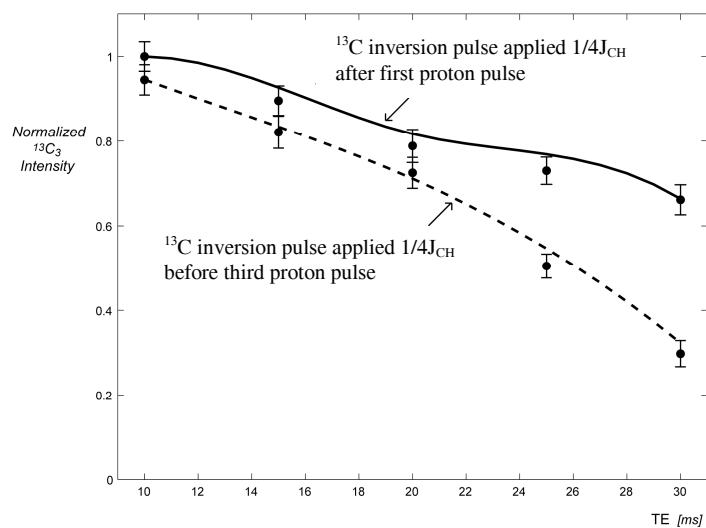


Figure 2: $^{13}\text{C}_3$ -Glu signal as a function of TE in response to the two versions of the sequence shown in Fig. 1(c). The experimental data confirm the numerical calculations. The error bars are half the peak-to-peak noise.