Accurate temperature measurements in the presence of field inhomogeneities

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Target audience: Clinical researchers involved in MR-guided thermal ablation work or, more generally, in field mapping.

Purpose: The goal of this work is to improve the measurement of temperature in the presence of B_0 inhomogeneities. Transcranial MRgFUS can be used to treat essential tremor^{1,2} near the thalamus. However, B_0 inhomogeneities near the treatment target often prevent accurate measurement³ of the therapy heating at the target site. With proton resonance frequency (PRF) thermometry, phase changes are converted into temperature changes by mapping the echo time, TE, to temperature. In the presence of field inhomogeneities, MR echoes may be translated and/or spread in k-space, potentially making the 'true' echo time vary spatially within the imaged object and take on values that differ from the nominal user-input TE value. Errors in TE lead to errors in temperature, and the purpose of this work is to describe and correct for these errors.

Methods: Consider a dual-echo gradient-echo sequence with imaging bandwidth BW_1 and BW_2 and nominal user-input echo time TE_1 and TE_2 for the two echoes, respectively. A relative temperature measurement could be obtained independently from each echo, and combined through a weighted sum with weights w_1 and w_2 : $\Delta T = (w_1 \Delta T_1 + w_2 \Delta T_2)/(w_1 + w_2)$. It can be derived, when selecting weights such that $w_2 = -w_1 \times BW_2/BW_1$, errors in ΔT_1 precisely cancel out errors in ΔT_2 , giving rise to a temperature measurement immune to field inhomogeneities. Ideally, however, weights should be chosen to optimize the temperature-to-noise-ratio (TNR), based on expected signal strength and/or TE value, and constraining the choice of these weights would have a TNR penalty. Alternately, inhomogeneity-induced errors in TE could be explicitly calculated, allowing any weights w_1 and w_2 to be subsequently selected for optimum TNR. With $\Delta \phi$ the measured temperature-induced phase shift, and τ the evolution time available for phase accumulation (in most sequences τ is simply equal to TE but in more exotic cases it may take on different and even negative values⁴), one can obtain:

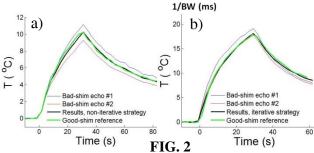
field inhomogeneities, making it difficult to perform pixel-by-pixel processing. A more general but more computationally-involved method is introduced below that remains valid even for a single-echo sequence or for a dual-echo sequence with $BW_1 = BW_2$.

Given a field map $B_0(\vec{r})$, spatial gradients in $B_0(\vec{r})$ are responsible for shifting signal in k-space.

A pulse-sequence specific function f_s is introduced that returns the time at which any given k-space point is sampled. The error in the echo time is given by:

$$\Delta TE_{j,i}(\vec{r}) = f_s(\Delta \vec{k}_{j,i}) - TE$$
, with $\Delta \vec{k}_{j,i}(\vec{r}) = 2\pi \times \gamma \times \hat{\tau}_{j,i} \times \nabla B_{0,i}(\vec{r})$, and $\hat{\tau}_{j,i} = TE_j + \Delta TE_{j,i-1}$,

where j is the echo number, i the iteration number, $\Delta TE_{j:0} = 0$. Corrected temperature measurements are obtained from corrected τ values. The field map also needs to be corrected in the process, hence the 'i' dependence in $B_{0,i}$. For example, with a dual-echo sequence, the field map would be obtained by comparing the phase of the two echoes and dividing by $(TE_{2,i} - TE_{1,i})$, and thus is corrected from one iteration to the next. A very small number of iteration (about 3) appeared to be sufficient in the present work.



Results: A series of thermometry experiments were performed by heating a gel

phantom using an ultrasound transducer inside a 3 T MR system (55 W for 20 s, 1.5×1.5×5 mm³, 128×128, *TE/TR*=10/17 ms, 30° flip angle). Temperature measurements were made in the presence of field inhomogeneities (introduced by intentionally de-adjusting shim parameters) as well as with carefully-shimmed scans and using a number of different settings for the sampling bandwidth, *BW*. Figure 1 shows that as expected the presence of field inhomogeneities cause errors in the measured temperature values. These errors take the form of a multiplicative factor that increases linearly with 1/*BW*, with a slope that depends on local field gradients at the transducer's focus. Through careful shimming, local field gradients could be removed, leading to accurate measurements at all sampled 1/*BW* settings ('o' symbols in Fig. 1). The proposed work aims at enabling accurate measurements despite field inhomogeneities.

Results in Fig. 2a were reconstructed with weights w_1 and w_2 chosen as described above ($BW_1 = -BW_2$, $w_1 = w_2$, 37 W for 30 s, $TE_1/TE_2/TR = 7.2/17.8/25$ ms), and reference results were obtained from carefully-shimmed scans. For Fig. 2b, a sequence with $BW_1 = BW_2$ was used instead, and data were reconstructed with the iterative method described above (57 W for 30 s, $TE_1/TE_2/TR = 6.2/16.7/24$ ms). In both cases please note that corrected results obtained in bad shim conditions, once reconstructed as proposed, were in good agreement with reference good-shim results.

Discussion: PRF thermometry measures temperature-induced field changes, and is in nature a field-mapping method. The type of errors described and corrected for here would presumably be present in any field-mapping MR application, not just thermometry. These errors can be quite significant, especially when a narrow imaging bandwidth is chosen to help boost SNR, and errors by as much as 26% were readily observed and corrected for.

Conclusion: Field inhomogeneities create errors in temperature measurements, which can be detected and corrected using the proposed strategies.

References: 1. Elias et al, N Engl J Med 369:7 (2013). 2. Lipsman et al, Lancet neurology. 12:5 (2013). 3. Sprinkhuizen, Chapter 7, Ph.D. thesis (2010). 4. Madore et al, MRM 66:658(2011). Support from grants R01CA149342, P41EB015898 and R01EB010195 is acknowledged.