## Simultaneous Determination of Electrical Properties and Proton Density in a Generalized Gradient-based Electrical **Properties Tomography**

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**Purpose** The original electrical properties tomography (EPT) was proposed based on a homogeneous Helmholtz equation to extract conductivity  $\sigma$ and permittivity  $\varepsilon$  (EP) of tissue noninvasively using MRI<sup>1,2</sup>. However, near boundaries separating different electrical properties, this simplified formula does not hold anymore, resulting in significant boundary artifacts<sup>3</sup>. In a previous study<sup>4</sup>, we have proposed a gradient-based EPT algorithm (gEPT), in which it was shown that the gradient g of EP can be obtained to provide EP maps with significantly improved boundary reconstruction and robustness against measurement noise. In this abstract, we introduce a more generalized gEPT framework to be able to derive g under realistic situations when the measured receive B1 field is biased by the unknown proton density  $(\rho)$ . Our results show that both EP and proton density  $\rho$  can be reliably estimated using this proposed approach.

**Theory** Based on the time-harmonic Maxwell's equations, ignoring the unknown  $B_z$  component, Theory Based on the time-harmonic Maxwell's equations, ignoring the unknown  $b_{\varepsilon}$  component, inhomogeneous Helmholtz equations in the form of transmit  $(B_1^+)$  or receive  $(B_1^-)$  B1 fields can be expressed as Eqs. 1 or  $2^{4.5}$ , respectively, in which the complex permittivity  $\varepsilon_c = \varepsilon - i \sigma / \omega$ ,  $\omega$  the Larmor angular frequency,  $\mu_0$  magnetic permeability of free space, gradient  $\mathbf{g} = \nabla \ln \varepsilon_c$ , and  $\hat{B}_1^-$  denotes the measurable proton-density-weighted receive B1  $(\rho B_1^-)$  with  $\gamma = 1/\rho$ . Equations 1 and 2 can further be extended into the magnitude  $|B_1^+|$ ,  $|\hat{B}_1^-|$  and phase  $\phi^{\pm}$  components of  $B_1^+$  and  $\nabla^2 (\gamma \hat{B}_1^-) = -\omega^2 \mu_0 \gamma \hat{B}_1^- \varepsilon_c + \nabla (\gamma \hat{B}_1^-)^T \begin{bmatrix} 1 & -i & 0 \\ i & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \mathbf{g}$  (2)

$$\nabla^{2} B_{1}^{+} = -\omega^{2} \mu_{0} B_{1}^{+} \varepsilon_{c} + (\nabla B_{1}^{+})^{T} \begin{bmatrix} 1 & i & 0 \\ -i & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \mathbf{g} \quad (1)$$

$$\begin{bmatrix} 2 (\mathbf{v} \hat{\mathbf{p}}^{-}) - \omega^{2} \mu_{0} \mathbf{v} \hat{\mathbf{p}}^{-} \mathbf{c} + \nabla (\mathbf{v} \hat{\mathbf{p}}^{-})^{T} \begin{bmatrix} 1 & -i & 0 \\ i & 1 & 0 \end{bmatrix} \mathbf{g} \quad (2)$$

 $B_1^-$  fields. Using a multi-channel transmit/receive array RF coil,  $|B_{1i}^+|$ ,  $|\hat{B}_{1k}^-|$ , relative phase

difference  $\Delta \phi_i^+ = \phi_i^+ - \phi_n^+$  and transceiver phase  $\phi_{nk} = \phi_n^+ + \phi_k^-$  can be measured<sup>6</sup>, where j and k denote the index of transmit and receive channels, respectively, and n a reference transmit channel. Knowing those information is sufficient to derive the unknown  $\mathbf{g}$ ,  $\nabla \ln \gamma$ ,  $\nabla \phi_n^+$ ,  $\sigma$  and  $\varepsilon$  based on Eqs. 1 and 2. Once g and  $\nabla \ln \gamma$  are obtained, maps of  $\sigma$ ,  $\varepsilon$  and  $\rho$  can be reconstructed from the gradients using the finite difference method.

Methods Simulation The proposed generalized gEPT algorithm was first evaluated utilizing finite-difference time-domain (FDTD) simulation. A sixteen-channel RF microstrip array coil<sup>7</sup> was rebuilt in the simulation software SEMCAD and loaded with an anatomically realistic human head model (Duke). Sixteen channels of  $B_1^+$  and  $B_1^-$  were computed with a spatial resolution of  $2x2x2 \text{ mm}^3$ . Experiment To validate the proposed method, a saline-gel phantom was built, including three different components of various concentration of NaCl. The corresponding conductivity and permittivity of the three solutions at 298MHz were measured using a dielectric probe as shown in Table 1. MRI experiment was performed on a 7T Siemens MAGNETOM scanner, equipped with 16-channel RF power amplifier and 32-channel receiver. The sixteen-channel RF microstrip array coil, modeled in the simulation study, was utilized to transmit RF energy and receive MR signal. For each channel,  $|B_1^+|$  and  $|\hat{B}_1^-|$  were obtained using a hybrid B1 mapping technique<sup>6,8</sup>. Relative phase  $\Delta \phi_j^+$  and transceiver phase  $\phi_{jk}$  were obtained from the signal phase of large flip angle gradient echo recalled (GRE) sequences. A  $\Delta B_0$  map was measured with GRE phase maps of two different echo times, to correct  $\phi_{jk}$  from

**Results** The reconstructed maps of conductivity  $\sigma$ , permittivity  $\varepsilon$  and proton density  $\rho$  of the simulation study are shown in Fig. 1 in comparison with corresponding target maps of the model input. Reconstructed relative errors (RE) and correlation coefficients (CC) are  $RE_{\sigma}$ =9.2%,  $RE_{\varepsilon}$ =9.8%,

 $RE_{\rho}$ =4.5%,  $CC_{\sigma}$ =0.98,  $CC_{\varepsilon}$ =0.85 and  $CC_{\rho}$ =0.97. Results of reconstructed  $\sigma$ ,  $\varepsilon$  and  $\rho$  of the phantom experiment are summarized in Table 1. Target  $\rho$  throughout the phantom is not directly measurable but assumed to be uniform (or normalized to be 1) because only NaCl concentration was slightly adjusted in the recipes. As we can see, using the proposed generalized gEPT approach, not only the electrical properties but also unknown  $\rho$  can be faithfully derived from the

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Table 1 Results of the phantom experiment. Units:  $\sigma - [Sm^{-1}]$ ;  $\varepsilon - [\varepsilon_0]$ ;  $\rho - [a.u.]$ .

		#1			#2			#3	
	$\sigma$	ε	ρ	σ	ε	ρ	σ	ε	ρ
Target	0.12	78		0.34	77		1.5	77	
Recon.	0.19	76	1.00	0.34	73	0.97	0.94	72	0.94
±Std.	±0.16	±5	±0.07	±0.06	±3	±0.03	±0.21	±2	±0.06
	-		-	•	•	-			

Figure 1 Results of simulation

measurable B1 field information. We do notice a significant drop of the reconstructed  $\sigma$  of phantom component #3. This could be due to errors of numerical differentiation near the vicinity of strong contrast of  $\sigma$  between #3 and #1. Finer imaging resolution may help improve the numerical performance.

Discussion and Conclusion In this study, we extended the generality of the previously proposed gEPT approach to quantitatively calculate gradient  ${f g}$  and unknown proton density ho at the same time. The method is generally applicable to any MR system with a multi-channel transmit/receive RF coil. Future work will be focused on advanced algorithms to combine the gradient  $\mathbf{g}$  and absolute  $\sigma$  and  $\varepsilon$ , the latter can also be derived in the solution but are more sensitive to noise.

Reference [1] Wen, Proc. SPIE 2003, 5030:471; [2] Katscher et al., IEEE Trans Med Imaging 2009, 28:1365; [3] Seo et al., IEEE Trans Med Imaging 2012, 31:430; [4] Liu et al., ISMRM 2013, 3486; [5] Zhang et al., IEEE Trans Med Imaging 2013, 32:1058; [6] Van de Moortele et al., ISMRM 2007, 1676; [7] Adriany et al. MRM 2008, 59:590; [8] Van de Moortele et al., ISMRM 2009, 367 Acknowledgement NIH R21 EB017069-01A1, R01 EB006433, R01 EB007920, R21 EB014353, R21 EB009138, T32 EB008389, P41 EB015894, 2R01 EB006835, 2R01 EB007327, S10 RR26783 and WM KECK Foundation.

inhomogeneous  $B_0$ .