

## Quantitative MR imaging method

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**Target Audience:** Researchers and clinicians interested in quantitative MR imaging.

**Purpose:** A patient receiving a typical MR exam can expect to lie in the scanner for as long as 45 min or so as the anatomy of interest gets repeatedly imaged, using a variety of tissue contrasts. The proposed approach allows the main MR parameters such as  $M_0$ ,  $T_1$ ,  $R_2$ ,  $R_2'$ ,  $B_1$  and  $B_0$  to be rapidly and quantitatively evaluated, so that any desired image contrast could, in principle at least, be computed rather than acquired.

**Methods:** Quantitative MR imaging performed within a practical scan time has been a quest for the MR community for a long time, with inspired contributions such as GESFIDE<sup>1</sup>, DESPOT<sup>2</sup>, MP-DESS<sup>3</sup>, TESS<sup>4</sup> and MR fingerprinting<sup>5</sup>. The method proposed here allows all of the main MR parameters to be calculated from data acquired from a single sequence. Previous methods may have targeted subsets of parameters for quantitative evaluation and/or required different sequences for different parameters. MR fingerprinting is a practical approach that also aims at evaluating all main parameters, but uses comparisons and matching with a databank of simulated results to do so. Because parameters are directly calculated here instead, a leaner approach that acquires only the minimum amount of data can be more naturally obtained.

The pulse sequence employed is a triple-pathway steady-state sequence similar to that used in TESS<sup>4</sup> (Fig. 1). Two scans, either sequential or interleaved, are performed with different nominal flip angles,  $\hat{\alpha}_1$  and  $\hat{\alpha}_2$ , and/or different TR settings,  $TR_1$  and  $TR_2$ .  $TE_{i,j,k}$  represents the echo time associated with the  $j^{\text{th}}$  echo for the  $k^{\text{th}}$  pathway, as sampled during the  $i^{\text{th}}$  scan. During TR the signal  $S_{i,k}(t)$  varies as described in Eq. 1<sup>6</sup>, where  $T_2 = 1/R_2$  and  $T_2^* = 1/(R_2 + R_2')$ , leading to the linear system in Eq. 2.

Using  $R_2$  and  $R_2'$  as found from Eq. 2, signals at  $t=0$  and  $t=TR$  can be calculated. The notation  $F_k$  and  $Z_k$  represents transverse and longitudinal magnetization states<sup>7</sup> and superscripts  $-$ ,  $+$  and  $\Rightarrow$  are used to distinguish between the moments just before, just after and a time  $TR$  after an RF pulse. In a sequential acquisition,  $Z_{i,k}^{\Rightarrow} = Z_{i,k}^-$ , while in an interleaved acquisition  $Z_{i,k}^{\Rightarrow} = Z_{h(i),k}^-$ , where the  $h(i)$  function returns the scan number that differs from  $i$  in a two-scan acquisition, i.e.,  $h(1) = 2$  and  $h(2) = 1$ . In the interest of space, only the sequential case is presented below. An important parameter referred to as 'mixing factor',  $X_{i,k}$ , is defined in Eq. 3. The flip angle  $\alpha_i$ , proportional to the  $B_1$  field, is found by solving Eq. 4, where  $c_i = \cos(\alpha_i)$ . To make Eq. 4 a single-unknown equation of  $\alpha_i$ , the relationship between  $c_2$  and  $\alpha_1$  must be known. In the small-flip angle and/or hard-pulse regime,  $c_2 = \cos(\alpha_1 \times \hat{\alpha}_2 / \hat{\alpha}_1)$ , while in the case of large slab-selective pulses simulations are used to evaluate  $g(\cdot)$  as in  $c_2 = g(\alpha_1, \hat{\alpha}_1, \hat{\alpha}_2)$ .

Once  $\alpha_1$  is found,  $T_1$  and  $M_0$  can be readily obtained from Eqs 5-6, where  $s_i = \sin(\alpha_i)$ . The presence of indices for  $T_1$  and  $M_0$  in Eqs 5-6 reflects the fact they can be obtained in independent ways, which are averaged to get a final value. In cases where more than one root can be found for Eq. 4, common sense points to the correct one:  $M_0$  and  $T_1$  are positive,  $\alpha_1 \approx \hat{\alpha}_1$  near the center of the slab, and the  $B_1$  field should be reasonably smooth.

The  $F$  and  $Z$  magnetization states are complex numbers, but great simplification can be derived from the realization that imperfections in the  $B_0$  and/or  $B_1$  fields tend to affect the phase of the measured  $F$  values but not their magnitude. By using real values for  $F$  as shown in Eq. 7, Eqs 3-6 can be solved using real rather than complex numbers. Phase information can be used, independently, to obtain a measure of the frequency offset.

**Results:** Simulating the effect of RF pulses and gradients on a virtual object led to simulated MR signals that were processed using the equations above. In the absence of noise, all parameters could be evaluated with essentially perfect accuracy:  $M_0$ ,  $T_1$ ,  $R_2$ ,  $R_2'$ , flip angle and frequency offset. The key parameter  $X_{i,k}$  is plotted in Fig. 2 for a range of  $\alpha$  and  $T_1$  values to provide a visual representation of how the method works ( $T_2 = 70$  ms and  $TR = 50$  ms were used). A single scan would not suffice to pinpoint both  $\alpha$  and  $T_1$  for a given value of  $X$  (e.g., blue lines). But two scans along with a known relation between them, for example  $\alpha_2 = 2 \times \alpha_1$ , can allow a solution to be found ('\*' marks in Fig. 2).

Phantom data have been acquired and processed as described above and results are shown in Fig. 4 ( $TR = 30$  ms, 4 echoes per TR per pathway,  $\hat{\alpha}_2 / \hat{\alpha}_1 = 10$ , matrix size =  $80 \times 80 \times 40$ , 192 s for 3D acquisition, selective Gaussian-shaped RF excitations). Nominal and independently measured values are shown in green. While some systematic errors could be readily noticed (e.g., ringing, linear trend for  $T_1$  along  $z$ ) very little noise could be seen in the results even though no filtering was applied.

**Discussion and conclusion:** The proposed approach allows  $M_0$ ,  $T_1$ ,  $R_2$ ,  $R_2'$ ,  $B_1$  and  $B_0$  to be evaluated with relatively short scan time over a 3D volume. At the time of writing, optimizations of acquisition parameters had not yet been performed. But as can be seen from the different curvature of lines above and below the  $180^\circ$  value in Fig. 2, considerable improvements in stability appear to be gained when working with greatly different  $\hat{\alpha}_1$  and  $\hat{\alpha}_2$  settings.

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$$|S_{i,k}(t)| \propto e^{-(R_2 + R_2') \cdot t}, k \geq 0; |S_{i,k}(t)| \propto e^{-(R_2 - R_2') \cdot t}, k < 0 \quad [1]$$

$$\ln(|S_{i,k}(TE_{i,j,k})|) = \begin{cases} \ln(|S_{i,k}(0)|) - TE_{i,j,k} \times (R_2 + R_2') & k \geq 0 \\ \ln(|S_{i,k}(0)|) - TE_{i,j,k} \times (R_2 - R_2') & k < 0 \end{cases} \quad [2]$$

$$X_{i,k} = 1 - 2 \times (F_{i,k}^- - F_{i,k}^+) / (F_{i,k}^- + F_{i,k}^+), |k| > 0 \quad [3]$$

$$(X_{1,k} c_1 - 1) \frac{TR_2}{TR_1} \times (X_{2,k} c_2 - 1) \times (X_{1,k} - c_1) \frac{TR_1}{TR_1} = 0, |k| > 0 \quad [4]$$

$$T_{1,i,k} = TR_i / \ln \left( \frac{X_{i,k} c_i - 1}{X_{i,k} - c_i} \right), |k| > 0 \quad [5]$$

$$M_{0,j,k} = \left( (F_{i,0}^+ - c_i \times F_{i,0}^-) + (F_{i,0}^- - c_i \times F_{i,0}^+) \times e^{\frac{TR_i}{T_1}} \right) / s_i \times \left( 1 - e^{\frac{TR_i}{T_1}} \right) \quad [6]$$

$$F_{i,k}^{+,-} = +|F_{i,k}^{+,-}|, k \geq 0; F_{i,k}^{+,-} = -|F_{i,k}^{+,-}|, k < 0 \quad [7]$$

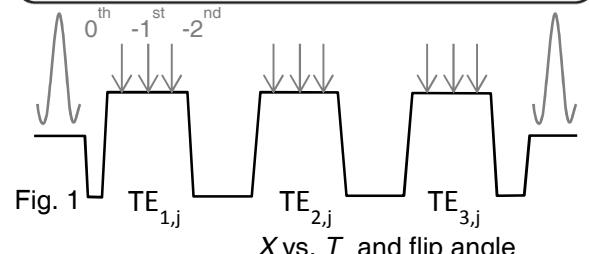


Fig. 1 X vs.  $T_1$  and flip angle

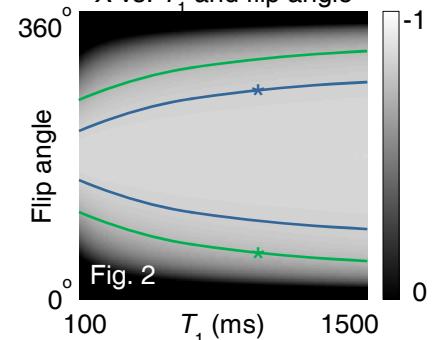


Fig. 2

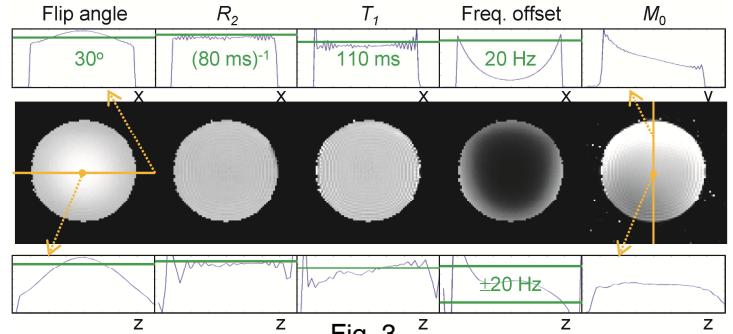


Fig. 3