

## B<sub>1</sub> Estimation using Adiabatic Refocusing: BEAR

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**Introduction:** Accurate measurement of B<sub>1</sub> transmit fields is important for calibration of the transmit system and quantitative MRI. We describe a novel phase-based B<sub>1</sub> estimation method using adiabatic refocusing (BEAR). Some important characteristics of the BEAR method are that the B<sub>1</sub> measurement is insensitive to off-resonance, T<sub>1</sub> and T<sub>2</sub>. BEAR also provides good image quality even in regions of B<sub>0</sub> inhomogeneity due to its robust spin-echo acquisition. We validate BEAR's performance in simulation and experimentally with comparison to Bloch-Siegert<sup>1</sup> (BS) B<sub>1</sub> measurements.

**Methods:** The second echo in a spin-echo sequence using two repeated adiabatic full passage (AFP) pulses will have no phase variation over the slice profile<sup>2</sup>. The BEAR method relies on the novel observation that by changing the relative magnitude of the two AFP pulses the phase of this echo will depend approximately linearly on B<sub>1</sub> and with very little variation over the slice profile. Fig. 1 shows the BEAR sequence with two sech<sup>3</sup> AFP pulses of magnitude  $\delta B_{1nom}$  and B<sub>1nom</sub>, where  $\delta$  is a scaling factor and B<sub>1nom</sub> is the nominal peak B<sub>1</sub> of the second AFP pulse. Numerical Bloch simulations were used to determine the signal phase dependence on B<sub>1</sub> for this sequence.

The sech pulses were designed with parameters T/β/μ equal to 12ms/800rad·s<sup>-1</sup>/5.5 giving a BW of 1.4kHz. The adiabatic threshold B<sub>1A</sub>, which we define as the minimum B<sub>1</sub> that ensures refocusing of 90% M<sub>xy</sub>, for this pulse is 0.095G. Assuming a B<sub>1nom</sub> of 0.175G, then  $\delta B_{1nom} > B_{1A}$  for  $\delta > 0.54$ . The BS method used an 8-ms Fermi pulse, with off-resonant frequency of ±4 kHz. A tip angle of 42°, TE of 44ms and TR of 500ms were used with a 2DFT acquisition on a GE Signa Excite 1.5-T scanner. To eliminate unwanted phase effects, phase-difference images were made from multiple acquisitions. For BEAR, the second acquisition reversed the order of the two adiabatic pulses; for BS, the second acquisition negated the off-resonant frequency of the Fermi pulse.

Imaging could be confined to a specified volume by making the refocusing pulses selective in Y (Fig. 1), and limiting the X readout receiver bandwidth. Fast, 1D projections could also be acquired using a single readout with k<sub>y</sub> = 0. For comparison to these fast projection acquisitions, 2D B<sub>1</sub> maps were also acquired, and their B<sub>1</sub> magnitude averaged along Y.

**Results:** Fig. 2a shows Bloch simulation results of BEAR's signal dependence on B<sub>1</sub> and  $\delta$ , with approximately linear phase dependence on B<sub>1</sub> for B<sub>1</sub> > B<sub>1A</sub>. The simulated magnitude and phase of the refocused M<sub>xy</sub>, as a function of B<sub>1</sub> and off-resonance frequency (Fig. 2b,c), illustrate BEAR's insensitivity to off-resonance over the effective bandwidth of the refocusing pulses. For  $\delta = 0.7$  and B<sub>1nom</sub> = 0.175G, the phase sensitivity was 80 rad/G, exceeding that of the BS method of 52 rad/G (Fig. 2a).

BEAR B<sub>1</sub> maps closely match BS B<sub>1</sub> maps (Fig. 3), with an average deviation from BS of 0.14% (phantom) and 1.5% (*in vivo*). Note, the BS method has B<sub>1</sub> map variations in areas of high B<sub>0</sub> inhomogeneity, causing increased deviation between the methods near the perimeter of the head. Scans repeated with a TR of 100ms showed similar results. Fig. 4 shows that B<sub>1</sub> projections acquired with BEAR are in agreement with projections of 2D B<sub>1</sub> magnitude maps, with less than 1.6% difference.

**Discussion and Conclusion:** The BEAR method is a novel method of B<sub>1</sub> mapping that can be localized to a slice or 3D block volume with a spin-echo acquisition that is appropriate for fast projection measurements. As the method measures transverse magnetization phase perturbation, it is insensitive to T<sub>1</sub> and T<sub>2</sub>. The method has a large dynamic range as long as the AFP pulses operate over their adiabatic threshold. Its sensitivity increases with increasing ratio (1/δ) of the refocusing pulse magnitudes. With the parameters used here, BEAR has sensitivity that is 153% of the BS method. However, the BEAR method has high SAR which can limit TR, and imposes a moderately long TE which can result in low signal for regions of short T<sub>2</sub>. Nevertheless, BEAR's high dynamic range, insensitivity to B<sub>0</sub>, T<sub>1</sub>, and T<sub>2</sub>, ability to make fast projection measurements, and linear quantitative relationship between phase and B<sub>1</sub> make it an ideal candidate for use in robust transmitter gain calibration.

**References:** [1] Sacolick *et al.*, MRM 63:1315–1322, 2010 [2] Conolly *et al.*, MRM 18:28–38, 1991 [3] Silver *et al.*, JMR 59:347–351, 1984 [4] Conolly *et al.*, JMR 83:549–564, 1989

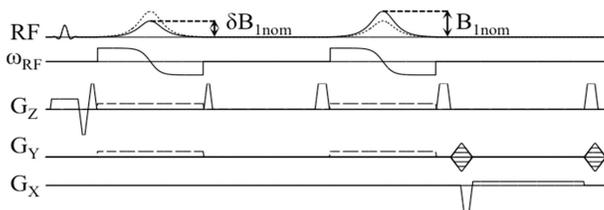


Figure 1: The BEAR sequence: two sech pulses generate a twice-refocused spin-echo. The refocusing pulses can be made selective in Z or Y for slice- or volumetric- imaging.

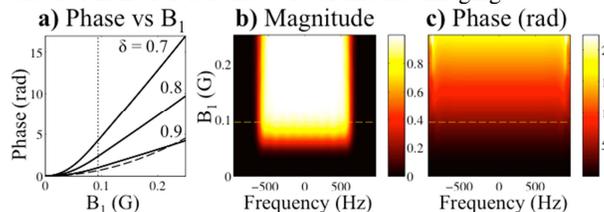


Figure 2: (a) BEAR's signal phase dependence on B<sub>1nom</sub> and  $\delta$  determined by Bloch simulations (solid). The BS phase dependence on B<sub>1</sub> for an 8-ms 4kHz offset Fermi pulse is shown for reference (dashed). (b) Magnitude and (c) phase plots for Bloch simulations of the slice profile for BEAR.

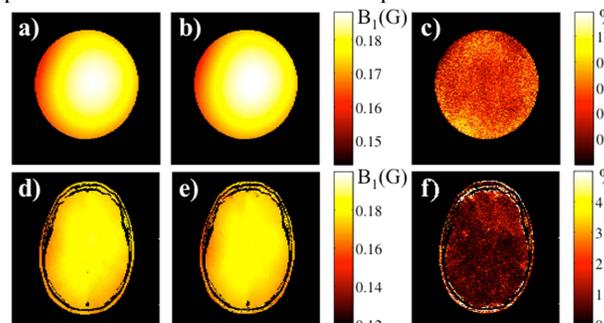


Figure 3: A comparison of B<sub>1</sub> maps: (a) BEAR phantom to (b) BS phantom, and (d) BEAR *in vivo* to (e) BS *in vivo*, with respective BEAR deviations from BS (c,f). Average deviations: 0.14% (c) and 1.5% (f).

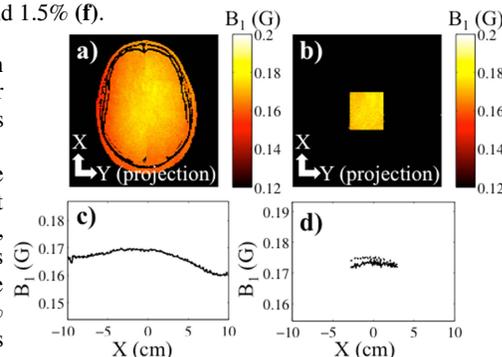


Figure 4: *In vivo* B<sub>1</sub> maps for: (a) slice and (b) volumetric scans. (c,d) B<sub>1</sub> projections (solid) and averages (dashed) of (a,c), with difference < 1.6%.