MEASUREMENT OF POST-MORTEM BRAIN MICROSTRUCTURE USING A CLINICAL MR SCANNER WITH OSCILLATING GRADIENTS

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Introduction Diffusion of water molecules in the presence of restricting boundaries reduces the apparent diffusion rates if the boundary geometry is of a similar order of magnitude to the diffusion distance. The ability to use this property to estimate the underlying restrictive geometry has long been recognized, for example, to yield estimates of axon diameter in white matter (1). While early approaches used standard gradient pairs, oscillating gradients at a particular frequency we have also been proposed to estimate pore geometry (2). For example, cosine- and sine-modulated oscillating gradients have been explored, and the former has been used to measure intracellular diameter (3). The theory underlying oscillating gradient experiments predicts that the diffusivity of completely restricted spaces approaches zero at $\omega = 0$, while hindered, but non-restricted, extracellular spaces have a finite diffusivity. We begin by considering simulations of these two compartments to understand their frequency dependence, which is poorly understood for the extracellular space. This motivates the combination of oscillating gradient waveforms and a modified signal model for estimating intra- and extra-cellular spaces. We present results in a post-mortem brain using a clinical scanner, suggesting the feasibility of in vivo measurement.

Methods Simulations: Camino (4) Monte Carlo simulations were used to generate diffusion-attenuated signals inside and outside of square-packed, non-abutting cylinders with impermeable walls. Simulations of cylinders with 10 and 3 μm radii with 20.6 and 6.18 μm separation, respectively, were performed. Cosine-modulated oscillating gradients with frequencies of 4-600 Hz were simulated perpendicular to the cylinder axis. The signals were converted into diffusivity D using $D = -1/b \ln(S/S_0)$ where b is the diffusion weighting and S and S_0 are the simulated signals with and without diffusion weighting, respectively. Eq. [1] is a modification of an existing analytic model (2) for the intracellular diffusivity spectrum. Eq. [1] was used to fit both the intracellular diffusivity spectra $D_1(\omega)$ and extracellular diffusivity spectra $D_{E}(\omega)$. B_{k} and a_{k} depend on geometry (as in Ref. 2) except that the radius in B_{k} and a_{k} now refers to an effective intracellular radius $R_{\rm I}$ or extracellular radius $R_{\rm E}$. $D_{\rm offset}$ is the diffusivity at $\omega = 0$ (zero for $D_{\rm I}(\omega)$). D_{∞} is the asymptote and equivalent to the free diffusivity of water D_{free} for $D_{\text{I}}(\omega)$. Experiment: Diffusion-weighted sequences with cosineand sine-like gradients (Fig. 1) were implemented on a 3 T MR scanner. Eight coronal slices centered on the splenium of a perfusion-fixed macaque brain were acquired using an 11" loop surface coil with FOV = 96×96 mm, thickness = 1.5 mm, TE/TR=150/1800 ms, 64×64 matrix, 63 averages, and four b=0 scans. Oscillating frequencies for cosine were 18.3, 20.0, 45.0, 50.1, 55.0, 60.2, 63.9, and 90.0 Hz and for sine, 18.3, 20.0, 85.0, 90.0, and 91.0 Hz. The maximum number of oscillations was limited by the TE. The diffusion attenuation in an ROI containing only the splenium was fit to

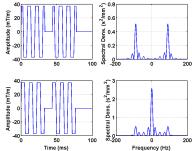


Fig. 1: Cosine- and sine-like diffusionweighting gradients (left) and the spectral densities of their integrals (right).

Eq. [2] and [3]. v_1 is the intracellular volume fraction, $b_{cos}(\omega)$ and $b_{sin}(\omega)$ are the diffusion weightings of the cosine- and sine-like diffusion gradients, and $D_{\rm I}(\omega)$ is the intracellular diffusivity spectrum from Ref. 2. $D_{\rm I}(\omega)$ includes the unknown intracellular effective radius R_1 and free diffusivity of water D_{free} , $D_{\text{F}}(\omega)$ includes the unknown extracellular effective radius R_{F} , D_{offset} , and D_{∞} as given in Eq. [1].

All fitting was done using Matlab.

Results Simulations: Both the intracellular extracellular and diffusivity spectra provided an excellent fit to the simulations (Fig. 2). For cylinders with 10 µm and 3 µm,

$$D(\omega) = D_{\text{offset}} + \sum_{k} B_{k} \frac{a_{k}(D_{\infty} - D_{\text{offset}})\omega^{2}}{a_{k}^{2}(D_{\infty} - D_{\text{offset}})^{2} + \omega^{2}}$$

$$E_{\cos}(\omega) = S_{\cos} / S_{0} = v_{I}e^{-b_{\cos}(\omega)D_{I}(\omega)} + (1 - v_{I})e^{-b_{\cos}(\omega)D_{E}(\omega)}$$

$$E_{\sin}(\omega) = S_{\sin} / S_{0} = v_{I}e^{-b_{\sin}(\omega)D_{I}(\omega)} + (1 - v_{I}) \left[e^{-b_{\sin}(\omega)D_{E}(\omega)} - e^{-b_{\sin}(\omega)D_{E}(\omega)} \right]$$
[3]

$$E_{\cos}(\omega) = S_{\cos} / S_0 = v_I e^{-b_{\cos}(\omega)D_I(\omega)} + (1 - v_I)e^{-b_{\cos}(\omega)D_E(\omega)}$$
[2]

$$E_{\sin}(\omega) = S_{\sin} / S_0 = v_I e^{-b_{\sin}(\omega)D_I(\omega)} + (1 - v_I) \left[e^{-b_{\sin}(0)D_E(0)} e^{-b_{\sin}(\omega)D_E(\omega)} \right]$$
[3]

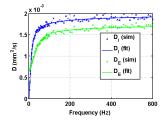


Fig. 2: Simulated and fitted diffusivity spectra for cylinders with $R_I = 10 \mu m$.

Eq. [1] estimated intracellular effective radii R_1 of 9.8 and 3.0 μ m, respectively, and extracellular effective radii R_E of 4.3 and 1.32 μm, respectively. Circles inscribed in the extracellular space would have radii of 4.6 and 1.37 μm, respectively, consistent with the fit values. Experiment: The fitting (Fig. 3) to Eq. [2] and [3] returned values of $v_1 =$ $0.60,\,R_1=9.5\,\,\mu m,\,R_E=5.4\,\,\mu m,\,D_{free}=0.65\times 10^{-3}\,\,mm^2/s,\,D_{offset}=3.7\times 10^{-3}\,\,mm^2/s,\,and\,D_{\infty}=8.6\times 10^{-3}\,\,mm^2/s.$ These values are in broad agreement with known intracellular volume fraction (0.7–0.8) and free water diffusion ($D = 2.0 \times$ 10^{-3} mm²/s), although the axon diameters are larger than expected from histology (~1 μ m).

Discussion We demonstrate the estimation of intracellular volume fraction, radius and effective radius of the extracellular compartment (related to packing density) using a novel set of cosine and sine oscillating gradient measurements. The fitted values for v_1 , R_1 , R_2 , and D_{free} are broadly consistent with those of fixed white matter measured with conventional techniques on high-performance pre-clinical scanners, as well as being consistent with our simulated diffusion spectra for the separate compartments. The diameter estimates are higher than expected based on histology, which is a general trend in reported measurements due to long diffusion times. D_{offset} likely relates to the tortuosity of the extracellular space and is similar to the tortuosity term in the intracellular model (2); however, we were unable to obtain reasonable fits of this model to our simulated extracellular signal curves. The oscillating gradients in our experiment utilize trapezoidal approximations to maximize the diffusion weighting while mainly preserving the frequency content. The gap in measurement frequencies in the sine-like gradients (Fig. 3) were due to spectral side lobes that make a single frequency assignment inappropriate. The fitted line does not contain any points in the gap since the

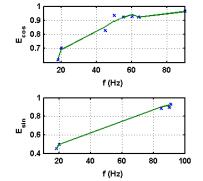


Fig. 3: Diffusion attenuation of measured and fitted values for cosine- (top) and sine-like (bottom) diffusion gradients.

fitting is dependent on the b-value at each measurement frequency. Future work will consider methods to attenuate these side lobes through careful design of the gradients and whether the combination of cosine and sine waveforms is critical.

Conclusions A model that fits the extracellular diffusivity spectrum of square-packed, non-abuttting cylinders was developed and validated with simulation. The combination of sine and cosine gradients provided sensitivity to the extracellular compartment at $\omega = 0$. Values for intracellular volume fraction, effective intracellular and extracellular radii, and diffusivity spectra of the splenium of a post-mortem brain were acquired using a clinical MR scanner.

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References: 1. Stanisz MRM 1997. 2. Stepišnik Physica B 2001. 3. Does MRM 2003. 4. Cook ISMRM 2006. 5. Barazany Brain 2009.