## Study of concomitant fields in multipolar PatLoc imaging

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## Introduction:

For spatial encoding in magnetic resonance imaging only the  $B_z$  component of the magnetic vector is of interest. In conventional imaging  $B_z$  is varied linearly along the three directions x, y and z to allow for orthogonal and rectangular voxels. One source of voxel distortion or signal dephasing for particular sequences are the concomitant fields  $B_x$  and  $B_y$  which are unavoidable to fulfil Laplace's equation. For linear gradients, e.g. for the x-gradient, the component  $B_z$  along x is transferred into  $B_x$  along z, Figure 1a/b. With the introduction of non-linear magnetic fields for imaging [1], concomitant fields of higher order harmonics become less trivial. This work studies the concomitant fields of higher order harmonics and of a simplified model of a PatLoc gradient coil [2].

## **Materials and Methods:**

Magnetic fields inside a source-free region can be calculated from Laplace's equation for the magnetic scalar potential  $\nabla^2\Psi=0$  [3]. Its solution is given by an infinite series of spherical harmonics [4]

$$\Psi = \sum_{m=0}^{\infty} \sum_{n=m}^{\infty} (A_n^m \cos(m\varphi) + B_n^m \sin(m\varphi)) r^n P_n^m (\cos\theta)$$
 (1)

where  $P_n^m$  are the associate Legendre polynomials and  $r,\theta,\phi$  are polar coordinates. From this potential the gradients  $B_x$ ,  $B_y$  and  $B_z$  of the magnetic field can easily be derived. The maximum of the calculated  $B_z$  component inside a unit sphere is used to normalise the magnetic fields to  $\pm 1$ . The total volume in which the magnetic fields are calculated is in arbitrary units of  $\pm 2$ , corresponding to a typical bore diameter to be twice the target spherical region. With the given  $B_x$  and  $B_y$  components the resulting concomitant field is calculated  $B_r = \sqrt{(B_x^2 + B_y^2)}$ .

In order to study the impact of the finite coil length, longitudinal and return current paths a step away from harmonic expansions towards wire patterns is required. Second order harmonics are of primary interest here, corresponding to the available PatLoc coil [3]. It is, however, difficult to draw general conclusions from a particular coil implementation. Therefore a vastly simplified model with a low number of parameters has been developed. This model includes a cylindrical magnet bore with a radius 2 with two flat single loop coil elements with a radius 2 placed atop the bore liner separated by a distance of 1/5 to represent one pole of the quadrupolar magnetic field. Total of 4x2 coil elements with alternating current directions are required to mimic a PatLoc gradient coil.  $B_x$ ,  $B_y$  and  $B_z$  were calculated based on an analytic field formula of a circular current loop. The simulated magnetic field components were scaled to normalize  $B_z$  to  $\pm 1$  inside the unit sphere.

## Results & Discussion:

For the linear x-gradient (A11) there is no radial component, which corresponds to the common assumption about the concomitant fields of linear gradients [5]. Figure 1a/b shows the scaled  $B_z$  component and the corresponding concomitant field  $B_x$ . Figure 2a/b displays the concomitant field  $B_r$  of the second order harmonics A22 (corresponding to the PatLoc encoding field) at z=0 and a 1D representation to visualize the distribution along y=0. Figure 3a/b displays the same plots for the concomitant field  $B_r$  of the third order harmonics. The 1D distributions suggest that the radial dependency of the concomitant field  $B_r$  is n-1 with the order of the harmonics. Figure 4a displays the encoding component  $B_z$  of the model PatLoc coil and Figure 4b the concomitant component  $B_r$ .

For linear gradients the magnitude of the main and concomitant fields is of the same order. For higher order magnetic fields the concomitant fields have a n-1 distribution with respect to the order of the harmonics. For the calculated harmonics of A22, the maximum  $B_{\rm r}$  inside the sphere at z=0.5 is 0.8 (Figure 4a) compared to the calculations from the simplified model with 1. This emphasises how important a realistic wire path is to estimate the concomitant fields. It is also important to know the concomitant fields to evaluate their effect on imaging and also for

Fig.1a z component of A11 (x-Fig.1b x component of A11 (xgradient) in ZX plane gradient) in ZX plane Fig 3a Bz component of A33 Fig 2a Bz component of A22 -0.5 2.5 Fig 3b radial component of A33 Fig 2b radial component of A22 A22r y=0 z=1Fig 3c 1D plot through Fig2a Fig 2c 1D plot through Fig2a B4C masked XY z=0.5 0.3

Fig 4b r component of model

prediction of peripheral nerve stimulation which is induced by  $\vec{B}$  and not only by  $B_z$ .

**References:** [1] Hennig et al, MAGMA 21(1-2):5-14 (2008); [2] Welz et al, ISMRM 2009, p.3073; [3] Romeo Hoult MRM 1,44-65(1984); [4] Hillenbrand et al, Applied Magnetic Resonance 29, 39-64(2005); [5] Bernstein et al., MRM 1998 39(2): 300-308

Fig 4a z component of model

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