## A simple model of gas exchange in the lung for hyperpolarized Xe129

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 $\underline{\textbf{Introduction}} \ \ \textbf{Hyperpolarized} \ \ \underline{\textbf{129}} \textbf{Xe} \ \ \textbf{is} \ \ \textbf{a} \ \ \textbf{unique} \ \ \textbf{contrast} \ \ \textbf{agent} \ \ \textbf{for} \ \ \textbf{pulmonary} \ \ \textbf{MRI} \ \ \textbf{because} \ \ \textbf{of} \ \ \textbf{its} \ \ \textbf{solubility} \ \ \textbf{in} \ \ \textbf{lung} \ \ \textbf{tissue} \ \ \textbf{and} \ \ \textbf{blood}. \ \ \textbf{Driven} \ \ \textbf{by} \ \ \textbf{diffusion}, \\ \textbf{like} \ \ \textbf{other} \ \ \textbf{blood} \ \ \textbf{gases} \ \ (\textbf{O}_2, \ \textbf{CO}_2) \ \ \textbf{in} \ \ \textbf{the} \ \ \textbf{lung}, \ \ \textbf{the} \ \ \textbf{dissolved} \ \ \textbf{xenon} \ \ (\textbf{DX}) \ \ \textbf{is} \ \ \textbf{under} \ \ \textbf{rapid} \ \ \textbf{exchange} \ \ \textbf{with} \ \ \textbf{the} \ \ \textbf{free} \ \ \textbf{xenon} \ \ \textbf{gas} \ \ \textbf{in} \ \ \textbf{the} \ \ \textbf{alveolar} \ \ \textbf{space}. \\ \ \ \textbf{The problem} \ \ \ \textbf{The problem} \ \ \ \textbf{The problem} \ \ \ \textbf{The problem} \ \ \textbf{The problem} \ \ \textbf{The problem} \ \ \ \textbf{$ Quantification of xenon exchange in the lung can therefore not only measure the microstructure, but the function of the lung. With the large chemical shifts of DX—197 ppm for xenon in tissue and blood plasma (TP) and 217 ppm for xenon in the human red blood cells (RBC) (Fig. 1)—one can saturate DX without disturbing the xenon gas and observe the uptake (1), or replenishment (2), of DX due to exchange using chemical shift saturation recovery (CSSR, (3)). We present here a simple model of gas exchange that interprets both TP- and RBC-xenon signals as functions of gas-exchange time, which allows simultaneously quantification of lung density, microstructure, function, and physiology.

Methods In order to study xenon exchange in different lung compartments we must separate the TP and RBC xenon peaks in the DX signal. While the TP xenon resonates consistently at 197 ppm, the frequency of the RBC xenon is species dependent: for rabbits (4) and mice (5) it is close to the TP xenon and is between 197 and 203 ppm. In this case the entire dissolved xenon signal in the real/imaginary channel can be fitted to a double-Lorentizan

in the frequency domain after phase correction:  $F(f) = \sum_{l=1,2} \frac{a_l}{(f-f_l)^2 + w_l^2}$  [1], where  $f_l$  are the peak positions and  $a_l$  and  $w_l$  (l=1,2) are fitting parameters (4); for many other species the RBC xenon (rats – 211 ppm (2), dogs – 212 ppm (1), humans – 217 ppm (6)) is farther apart from the TP xenon and they cannot be fitted with a common phase factor due to the dephasing between them; instead, they can be fitted using the double-Lorentzian with individual

phase factors (
$$\theta_l$$
):  $F(f) = \sum_{i=1,2} \frac{a_i \exp(-i\theta_i)}{w_i + i(f - f_i)}$  [2], from which  $\text{Re}[F(f)] = \sum_{i=1,2} a_i \frac{w_i \cos\theta_i - (f - f_i)\sin\theta_i}{(f - f_i)^2 + w_i^2}$  [3] and  $\text{Im}[F(f)] = \sum_{i=1,2} -a_i \frac{w_i \sin\theta_i + (f - f_i)\cos\theta_i}{(f - f_i)^2 + w_i^2}$  [4] are the real and imaginary parts, respectively. The signal intensity of each peak can be obtained by numerically integrating the corresponding Lorentzian.

Theory The calculation is based on the solution of the 1D diffusion equation for DX in a simplified model of lung septum as shown in Fig. 1, where we assumed the total septal thickness and the air-blood barrier thickness to be d and  $\delta$ , respectively. As results, the normalized TP xenon signal  $S_{TP}(t)$  and the RBC xenon signal  $S_{RBC}(t)$  by the corresponding gas signal amplitudes as functions of the gas exchange time t are, respectively (7).

Explore signal S<sub>RBC</sub>(t) by the corresponding gas signal amplitudes as functions of the gas exchange time trace, respectively (7),
$$S_{\text{TP}}(t) = b \left[ \frac{2\delta}{d} - \frac{8}{\pi^2} \sum_{n=\text{odd}} \frac{1}{n^2} \left( 1 - \cos \frac{n\pi\delta}{d} \right) e^{-n^2t/T} \right] + b \left( 1 - \eta \right) \left\{ 2 \left[ \left( 1 - \frac{2\delta}{d} \right) \frac{t}{t_x} - \frac{8T}{\pi^2 t_x} \sum_{n=\text{odd}} \frac{1}{n^4} \cos \left( \frac{n\pi\delta}{d} \right) \left( 1 - e^{-n^2t/T} \right) \right] + \left( 1 - \frac{t}{t_x} \right) \left[ \left( 1 - \frac{2\delta}{d} \right) - \frac{8}{\pi^2} \sum_{n=\text{odd}} \frac{1}{n^2} \cos \left( \frac{n\pi\delta}{d} \right) e^{-n^2t/T} \right] \right\}, \quad [5]$$

$$S_{\text{RBC}}(t) = b\eta \left\{ 2 \left[ \left( 1 - \frac{2\delta}{d} \right) \frac{t}{t_{\chi}} - \frac{8T}{\pi^2 t_{\chi}} \sum_{n = \text{odd}} \frac{1}{n^4} \cos \left( \frac{n\pi\delta}{d} \right) \left( 1 - e^{-n^2 t/T} \right) \right] + \left( 1 - \frac{t}{t_{\chi}} \right) \left[ \left( 1 - \frac{2\delta}{d} \right) - \frac{8}{\pi^2} \sum_{n = \text{odd}} \frac{1}{n^2} \cos \left( \frac{n\pi\delta}{d} \right) e^{-n^2 t/T} \right] \right\}.$$
 [6]

There are 5 shared parameters between  $S_{TP}(t)$  and  $S_{RBC}(t)$ : b,  $\delta/d$ , T,  $\eta$  and  $t_X$ , where  $b = \lambda dS_A/V_g$  is the normalization constant and  $T = d^2/(\pi^2 D)$  is the gas-exchange time constant for xenon in the lung. The other parameters are:  $\lambda$  — average Ostwald solubility of xenon in tissue and blood;  $S_A/V_g$  surface-area-to-volume ratio (SVR); D — average diffusion coefficient of xenon dissolved in tissue and blood;  $\eta$  — fraction of RBC xenon in the blood;  $t_x$ — capillary transit time.  $t_x$  is introduced to take into account the effect of blood flow, which is illustrated in Fig. 2. From  $\eta$  the hematocrit (Hct) can be

estimated using  $\text{Hct} = \frac{\eta/\lambda_{\text{RBC}}}{\eta/\lambda_{\text{RBC}} + (1-\eta)/\lambda_{\text{P}}}$ , where  $\lambda_{\text{RBC}}$  and  $\lambda_{\text{P}}$  are the xenon solubilities in RBC and plasma, respectively.

Results Using the nominal parameter values for human ( $S_A/V_g$ =250 cm<sup>-1</sup>,  $\lambda$ ≈0.12, d≈10 μm,  $\delta$ ≈2 μm, T≈30 ms,  $\eta$ ≈0.5,  $t_X$ =1.5 s (8-12))  $S_{TP}(t)$  and  $S_{RBC}(t)$  in Eqs. [5] and [6], respectively, are plotted in Fig. 3. These curves in Fig. 3 qualitatively interpret the plots of data available in the literature (2,13-15). Discussion and Conclusions The validity of the proposed model of exchange can also be verified by examining its approximations. Specifically, adding  $S_{TP}(t)$  and  $S_{RBC}(t)$ , and setting  $\delta$ =0, we recover the expression for the entire DX signal (F(t)) when the barrier is ignored, which was first derived by Patz et al. (16); further, by setting 1/t<sub>X</sub>=0 we arrive at the simple result assuming no flow ([5] in (17)). In summary, we have presented a simple model of gas exchange which expresses the DX signals in terms of important pulmonary parameters. By fitting the xenon uptake data using the proposed model one can simultaneously measure lung microstructure, physiology, and function. It can potentially be used to detect a variety of pulmonary diseases.

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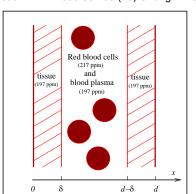


Figure 1 A simplified model of lung septum. Septal thickness = d, barrier thickness =  $\delta$ . Xenon in tissue and blood plasma resonates at 197 ppm, xenon in the red blood cells resonates at 217 ppm.

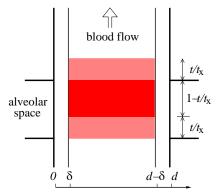
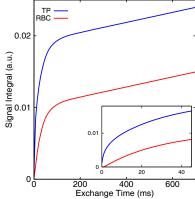


Figure 2 Effect of blood flow on gas exchange in the calculation of  $S_{RBC}(t)$ . The pink areas are partial volumes of exchange due to blood flow.



**Figure 3** Plots of  $S_{TP}(t)$  (blue) and  $S_{RBC}(t)$  (red) in Eqs. [5] and [6], respectively.