

"M_Y way" – a new construction technique for broadband slice-selective refocusing pulses

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Broadband spin-echo refocusing pulses are of interest not only for their highly nonlinear spin physics but also for their usefulness in PRESS spectroscopy, especially at higher field strengths where chemical shift misregistration can be a problem. A variety of techniques have been used to design these pulses, notably iterative minimization of a cost function [1-3], the Shinnar-Le Roux (SLR) algorithm [4], and SLR combined with subsequent flipping of **A** polynomial roots across the unit circle in a quest for a reduced B_{1max} [5]. The last approach has been used to create an entire series of phase-modulated (PM) pulses, but the process is slowed by the need to search for an optimal combination of "flippers" amidst a tangle of roots.

Yet another method is to first generate a slice-selective self-refocused 90° excitation pulse, namely one that needs no subsequent gradient rephasing. Combining this pulse with a time-reversed version of itself creates a 180° refocusing pulse [6]. However, the resulting state of out-of-slice magnetization halfway through – M_Z ≅ 1, M_X ≅ 0, M_Y ≅ 0 – is overly restrictive. For symmetric, purely amplitude-modulated (AM) refocusing pulses developed with other methods, we have found that the key magnetization characteristic halfway through the pulse is that M_Y ≅ 1 inside the slice and M_Y ≅ 0 outside. As such, rather than combining "traditional" self-refocused 90° pulses, one can specify merely the desired M_Y response ... and let M_X and M_Z run free.

The same approach can be used for higher-bandwidth phase-modulated pulses as well: start with either a slice-selective self-refocused PM 90° pulse or (*more generally*) a self-refocused PM 90° pulse with the desired M_Y properties. When combining these pulses with their time-reversed counterparts, however, the imaginary component of the second piece must be inverted [6]. This creates an amplitude-symmetric, phase-antisymmetric refocusing pulse with a simple relationship between the refocusing profile **E**(Δν) and the M_Z response profile: **E** = (1-M_Z)/2 (i.e., consistent pancake flipping about one axis in the x-y plane) [2,5]. (Note, however, that this symmetry property is useful but not essential for pulse design. Refocusing pulses with different symmetries – asymmetric AM or fully symmetric PM, for example – can be generated by iterative techniques if the refocusing profile itself is used as the in-slice target function [2].) Finally the freshly forged combination typically benefits from some additional fine-tuning and optimization as an actual refocusing pulse.

A variety of both AM and PM refocusing pulses have been created with the "M_Y way" technique, and some have found use in PRESS spectroscopy sequences at 3T.

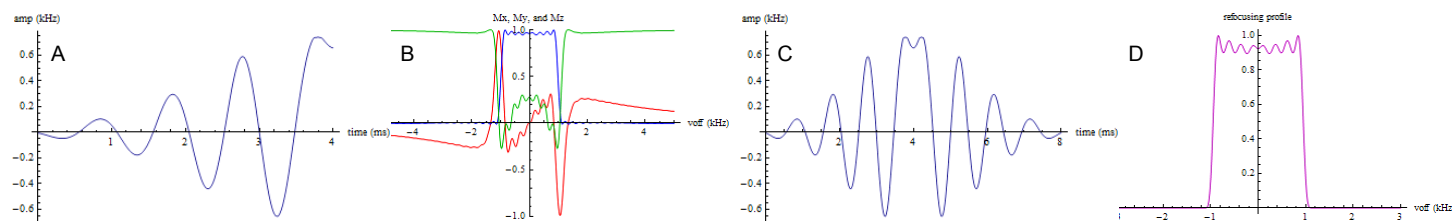


Figure 1. (A) A self-refocused "M_Y-design" pulse; (B) its M_X, M_Y, and M_Z profiles in red, blue, and green; (C) a symmetric pulse constructed from (A); (D) the refocusing profile of pulse (C).

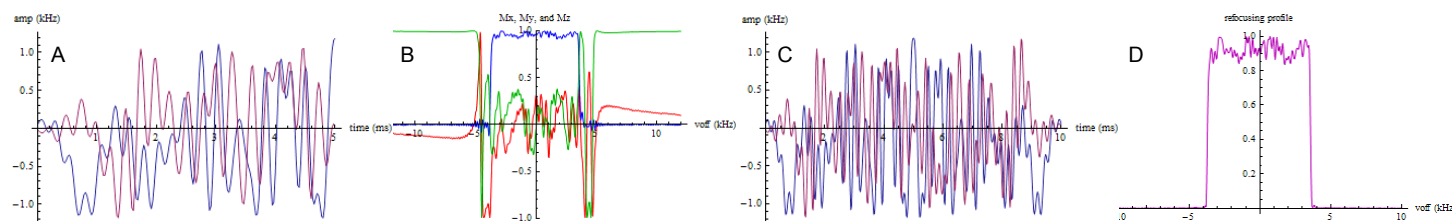
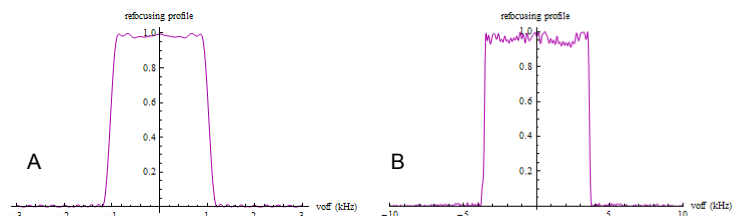


Figure 2. (A) A phase-modulated self-refocused "M_Y-design" pulse (*real and imaginary components displayed*); (B) its M_X, M_Y, and M_Z profiles in red, blue, and green; (C) an antisymmetric pulse constructed from (A); (D) the high-bandwidth refocusing profile of pulse (C).

Figure 3. Refocusing profiles for "refined" versions of the pulses displayed in Figs. 1 and 2. (Additional optimization was performed on the fused pulse halves to polish up the refocusing performance.) (A) The AM pulse from Fig. 1; (B) the PM pulse from Fig. 2.



References

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