# High-Resolution MR Spectroscopy in Inhomogeneous and Unstable Fields via Intermolecular Zero-Quantum Coherences

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## Introduction

Intermolecular zero-quantum coherences (iZQCs) are not susceptible to magnetic field inhomogeneities over distances much larger than the dipolar correlation distance, and can be used to achieve high-resolution spectra in inhomogeneous fields<sup>1</sup>. However, due to its intrinsic low SNR, iMQC spectroscopy requires more scanning time than its 1D counterparts. As a result, motions in *in vivo* studies will cause strong  $t_1$  noises in the 2D spectra of iMQCs and thus restrict previous attempt of high-resolution<sup>2</sup>. IZQCs have been utilized to achieve high-resolution NMR in unstable fields of a 25 T electromagnet<sup>3</sup>. In current work, the detailed mechanism to achieve high-resolution spectroscopy in unstable fields via iZQCs is presented and a modified iZQC pulse sequence, the iDQF-HOMOGENIZED<sup>4</sup> with stroboscopic acquisition (dubbed as SB-iDH), is designed to suppress inhomogeneous broadenings and motion artifacts in MRS.

# Methods

The desired evolution of a typical iZQC sequence can be given as:

$$I_z S_z \xrightarrow{\frac{\pi}{2}I_x, \frac{\pi}{2}S_x} \xrightarrow{\frac{\pi}{2}I_x, \frac{\pi}{2}S_x} \xrightarrow{\frac{\pi}{2}I_x} \frac{1}{4} I^- S^+(t_1) \xrightarrow{\frac{\pi}{2}I_x} \xrightarrow{\frac{\pi}{2}I_x} \frac{1}{4} I_z S^+ \xrightarrow{\pi I_x, \pi S_x, D_{IS}I_z, S_zt_2} \xrightarrow{\frac{\pi}{2}I_x, \frac{\pi}{2}S_x} \xrightarrow{\frac{\pi}{2}I_x} \xrightarrow{\frac{\pi}{2}I_x$$

where I spin (corresponding to solvent) and S spins (corresponding to solutes) are spin-1/2 systems. The iZQCs are not sensitive to static field inhomogeneity:

$$M_{s} \sim e^{i\omega_{l2QC}t_{1}}e^{i\omega_{sQC}t_{2}} = e^{i[-\omega_{l} - \Delta\omega(\mathbf{z}) + \omega_{s} + \Delta\omega(\mathbf{z})]t_{1}}e^{i[\omega_{s} + \Delta\omega(\mathbf{z})]t_{2}} = e^{i(-\omega_{l} + \omega_{s})t_{1}}e^{i[\omega_{s} + \Delta\omega(\mathbf{z})]t_{2}}$$

where  $\Delta \omega(\mathbf{z})$  denotes the angular frequency deviation at the position  $\mathbf{z}$ . However, it is more complicated in the case of unstable fields. The magnetic field varies from scan to scan as  $t_1$  increases:

$$M_{S} \sim e^{i\omega_{IZQC}t_{1}} e^{i\omega_{SQC}t_{2}} = e^{i[-\omega_{I} - \Delta\omega(\mathbf{z}, t_{1}) + \omega_{S} + \Delta\omega(\mathbf{z}, t_{1})]t_{1}} e^{i[\omega_{S} + \Delta\omega(\mathbf{z}, t_{1}, t_{2})]t_{2}} = e^{i(-\omega_{I} + \omega_{S})t_{1}} e^{i[\omega_{S} + \Delta\omega(\mathbf{z}, t_{1}, t_{2})]t_{2}}.$$
(3)

Although the unstable factor is eliminated in the iZQC term, the  $t_1$  dependence of  $\Delta \omega$  in the  $t_2$  acquisition will lead to strong  $t_1$  noises. From Eq. (3) it can be seen that the elimination of the chemical shift term in the  $t_2$  period can suppress the influence of  $t_1$ -dependent  $\Delta \omega$  without affecting the spectral information in the F1 axis. A CPMG scheme with stroboscopic acquisition is utilized for this attempt: one-point acquisitions take place at spin echo maxima between CPMG trains. As a result, chemical shift is refocused in the F2 dimension. The intermolecular double-quantum filter (iDQF) scheme<sup>4</sup> is also utilized for solvent suppression.

Experiment was performed on a Varian Unity+ 500 MHz NMR spectrometer. To simulate the inhomogeneous and unstable field environment, the shimming coils are deliberately detuned to produce a field inhomogeneity of 60 Hz, and background field gradients randomly vary from scan to scan in the range of -0.01 Gauss  $\cdot$  cm<sup>-1</sup> ~ 0.01 Gauss  $\cdot$  cm<sup>-1</sup>. A solution of methyl ethyl ketone (MEK) dissolved in cyclohexane is used. The parameters of coherence selection gradients (CSGs) are G' = 4 Gauss  $\cdot$  cm<sup>-1</sup> × 0.5 ms, G = 10 Gauss  $\cdot$  cm<sup>-1</sup> × 0.5 ms.  $100 \times 512$  points were acquired with spectral widths of  $100 \times 1500$  Hz. A 2-step phase cycling was used: the phases for the 2nd RF pulse and the receiver were (*x*, *y*) and (*x*, *-x*), respectively.

### Results

The results are presented in Fig. 2. It can be seen that high-resolution spectral information such as chemical shifts and *J*-coupling multiplets can be obtained in the SB-iDH spectrum. For comparison, an experiment with the same sequence and parameters but without stroboscopic acquisition was performed. Strong  $t_1$  noises can be observed in the 2D spectrum and severe artifacts in the projection conceal the "real" resonances. *In vivo* studies utilizing the SB-iDH sequence are under work.

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1. Vathyam S., et al., Science 272, 92, 1996.

- 2. Galiana G., et al., J. Am. Chem. Soc. 127, 17574, 2005.
- 3. Lin Y. Y., Phys. Rev. Lett. 85, 3732, 2000.
- 4. Chen X., et al., Phys. Chem. Chem. Phys. DOI: 10.1039/b709154k, 2007.



δ

Fig.1. The SB-iDH pulse sequence.

2G

(2)

(1)

Fig. 2. 'H NMR spectra of MEK in cyclohexane. (a) Conventional 1D spectrum in a homogeneous field, (b) the iZQC spectrum without stroboscopic acquisition in an inhomogeneous and unstable field, and (c) the iZQC spectrum obtained by the SB-iDH sequence in the same field condition as (b).