

# On the effects of a varied diffusion time in vivo: Is the diffusion in white matter restricted?

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**Introduction.** The signal-versus- $b$  curve from cerebral white matter (WM) can be modelled by a bi-exponential function, including a fast and a slow diffusion component [1]. Pulsed gradient stimulated echo (PGSTE) NMR experiments with varying the diffusion time ( $T_d$ ) have indicated that the slow diffusion component perpendicular to WM fibre bundles in excised nerves is restricted [2, 3]. However, the interpretation of such components in the context of *in vivo* MRI remains unclear [4, 5]. In previous MRI studies pulsed gradient spin echo sequences were typically employed, limiting the maximal  $T_d$  and the diffusion-measurement direction were, in contrast to NMR experiments, uncorrelated to the actual direction of the nerve fibre bundle (cf. [1]). In this study, measurements perpendicular and parallel to the corticospinal tract (CST) *in vivo* were performed using a PGSTE sequence to investigate whether effects of restricted diffusion could be observed as  $T_d$  was prolonged.

**Method.** Measurements were performed on 7 volunteers using a Siemens 3T Allegra unit. Initially, diffusion tensor (DT) imaging was executed in 12 diffusion encoding directions. A region of interest (ROI) in one slice containing the CST was selected and based on the eigenvectors of the DT in the ROI, new images were acquired with the diffusion encoding direction being either perpendicular ( $\mathbf{n}_\perp$ ) or parallel ( $\mathbf{n}_\parallel$ ) to the CST. Signal-versus- $b$  curves were acquired for  $T_d=64,100,144,196$  and 256 ms by varying the mixing time with TR/TE=2500/134 ms and  $\delta=50$  ms. For each  $T_d$ , 29  $b$ -values were sampled with  $b_{\max}=28\,000$  s/mm<sup>2</sup> along  $\mathbf{n}_\perp$  and 30  $b$ -values with  $b_{\max}=6\,000$  s/mm<sup>2</sup> along  $\mathbf{n}_\parallel$ . In an additional measurement, the signal-versus- $b$  curve was sampled using 87  $b$ -values up to 86 000 s/mm<sup>2</sup> ( $T_d=256$  ms) in  $\mathbf{n}_\perp$ . Bi-exponential and mono-exponential functions were fitted to the signal-versus- $b$  curves, and were considered acceptable if the index of satisfaction ( $IS$ ) was below 0.001 [1]. Furthermore, two-dimensional Monte Carlo simulations of  $5 \times 10^4$  molecules placed inside and between a lattice of cylindrical cells (diameter  $d$ ) were performed with sequence parameters as in the perpendicular measurements and with  $D_{\text{intra}}=D_{\text{extra}}=1.7$   $\mu\text{m}^2/\text{ms}$  and an intra-cellular fraction  $p_{\text{intra}} = 80\%$  for  $d=3,7,11$  and 15  $\mu\text{m}$ . In each step, the positions were updated with equal probabilities of a positive and a negative step ( $\Delta x=0.2$   $\mu\text{m}$ ). The membrane transition probability was given by  $p_0=m/n \cdot \Delta t/\tau_{\text{intra}}$ , where  $m$  is the number of possible displacements inside the cylinder and  $n$  is the number of displacements leading to a possible membrane transition and  $\Delta t=\Delta x^2/2D$ . For each  $d$ , five exchange times  $\tau_{\text{intra}}$  were simulated ( $\tau_{\text{intra}}=25,100,200,500$  and 1000 ms). Violation of the short gradient pulse (SGP) condition implies that expected diameters ( $d_{\text{exp}}$ ) are lower than the simulated  $d$  ( $d_{\text{exp}}$  were estimated from Fig. 2 in [6]). Finally, a two-compartment modified Kärger model, similar to the model in [7], was fitted to the data from the perpendicular direction as well as to the simulated signal values.

**Results.** The normalized signal perpendicular to the CST did not vary with  $T_d$  to any observable extent (Fig. 1), in contrast to the signal from the simulations (Fig. 2), where large effects on the signal-versus- $b$  curve were observed for varied  $T_d$ . In the measurements perpendicular to the CST, a bi-exponential function was required to fit the data well, but parallel to the tract a mono-exponential model was sufficient. The estimated parameters of the bi-exponential and mono-exponential fit

(Table 1) did not change significantly as  $T_d$  was varied, but  $ADC_{\text{fast}}$  and  $ADC_{\text{slow}}$  decreased and  $p_{\text{fast}}$  increased for the measurement with higher  $b_{\max}$ . SNR-values for  $b_{\min}$  and  $b_{\max}$  are also displayed in Table 1. The two-compartment model was found to fit well to the simulated signal values. Regarding the estimated parameters (Table 2), the estimated  $d$  values were in the expected range, except for  $d=3$   $\mu\text{m}$ . The estimated  $\tau_{\text{intra}}$  were all in good agreement with the simulated, but in general,  $p_{\text{intra}}$  was underestimated. Finally, the two-compartment model did not fit well to the *in vivo* signal values and no set of simulated signal versus  $b$ -curves resembled those obtained *in vivo*.

**Discussion.** In the present study,  $T_d$  was varied in approximately the same range as in previous NMR spectrometry studies [2,3], but in contrast to the NMR studies no obvious evidence of restricted diffusion was observed. For example, no increase of the normalized signal value was observed perpendicular to the CST as  $T_d$  was prolonged (Fig. 1) and the parameters of the bi-exponential fit did not change as  $T_d$  was varied (Table 1). The decrease in  $ADC_{\text{fast}}$  and  $ADC_{\text{slow}}$  as  $b_{\max}$  was increased (Table 1) shows that comparisons of parameters from a bi-exponential fit between measurements with different  $T_d$  should only be compared under conditions with identical  $b_{\max}$ . Since no effects of restricted diffusion perpendicular to the CST was observed *in vivo*, in contrast to similar *ex vivo* NMR studies, we hypothesise that differences between living and excised tissue might account for the observed differences.

The two-compartment exchange model was likely to be able to extract  $d$  and  $\tau_{\text{intra}}$  in the expected range when the diffusion occurs in cylindrical cells, but since no set of signal-versus- $b$  curves from the simulations resembled the signal curves obtained *in vivo*. Furthermore, the model did not fit well to the obtained signal values *in vivo*. Hence, our conclusion is that the diffusion perpendicular to WM tracts *in vivo* should be modelled in another manner than by the two-compartment Kärger model, for example as two slowly exchanging diffusion pools [5].

**References.** [1] Clark AC, *et al.* MRM 2000;44:852-859 [2] Assaf Y, Cohen Y. MRM 2000;43:191-199 [3] Nossin-Manor *et al.* MRM 2005;54:96-104 [4] Schwarcz A, *et al.* MRM 2004;51:278-285 [5] Le Bihan D. Phys Med Biol 2007;52:R57-R90 [6] Lätt J, *et al.* IEEE Trans Med Imag 2007;26 [7] Pfeuffer J, *et al.* NMR Biomed 1998;11:19-31

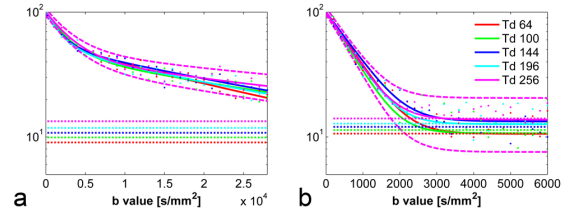


Figure 1. Normalized signal curves from a normal volunteer (a) perpendicular and (b) parallel to the CST. The fitted functions are displayed as solid lines. Dashed lines indicates one standard deviation of the signal for  $T_d=256$  ms. Note that the mean background noise level (dotted lines) appears to increase for increasing  $T_d$  due to the normalization of the signal.

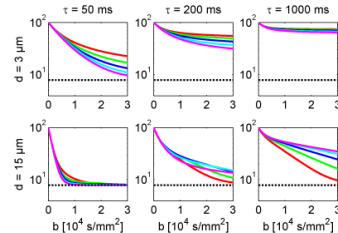


Figure 2. Example of the signal-versus- $b$  curves from the simulations. Dashed lines indicated the noise floor. Colour encoding as in Fig. 1.

Table 1. The parameters of the bi-exponential and mono-exponential fit perpendicular and parallel to the CST, respectively with  $T_d$  in ms,  $ADC$ :s in  $\mu\text{m}^2/\text{ms}$ ,  $p_{\text{fast}}$  in percent and SNR-values for  $b_{\min}$  and  $b_{\max}$  respectively. The bottom row ( $T_d=256^*$  ms) shows parameters from the measurement with higher  $b_{\max}$ . None of the parameters showed a significant change with  $T_d$  in a dependent ANOVA-test with  $p = 0.30, 0.22$  and  $0.25$  for  $ADC_{\text{fast}}$ ,  $ADC_{\text{slow}}$  and  $p_{\text{fast}}$  in the perpendicular direction and  $p = 0.42$  for  $ADC$  in the parallel direction ( $n=7$ ).

$T_d$	Perpendicular				Parallel	
	$ADC_{\text{fast}}$	$ADC_{\text{slow}}$	$p_{\text{fast}}$	SNR	$ADC$	SNR
64	$0.45 \pm 0.13$	$0.031 \pm 0.006$	$49 \pm 4$	16/3.8	$1.13 \pm 0.14$	12/1.4
100	$0.39 \pm 0.06$	$0.027 \pm 0.001$	$51 \pm 3$	14/3.7	$1.09 \pm 0.07$	11/1.4
144	$0.40 \pm 0.07$	$0.029 \pm 0.004$	$49 \pm 3$	13/3.4	$1.09 \pm 0.12$	11/1.4
196	$0.34 \pm 0.10$	$0.023 \pm 0.009$	$54 \pm 7$	12/3.1	$1.06 \pm 0.05$	10/1.3
256	$0.40 \pm 0.09$	$0.029 \pm 0.007$	$49 \pm 6$	11/2.9	$1.05 \pm 0.09$	9.2/1.4
256*	$0.22 \pm 0.06$	$0.014 \pm 0.004$	$62 \pm 5$	11/1.6		

Table 2. The simulated parameters along with the parameters estimated by the two-compartment Kärger model. The  $\tau_{\text{intra}}$ :s are averaged over the different diameters ( $n=5$ ) and the  $d$ :s over the different  $\tau_{\text{intra}}$ :s ( $n=4$ ). The expected  $d_{\text{exp}}$  was estimated from Fig. 2 in [6]. The  $p_{\text{intra}}$  was averaged over all simulated settings ( $n=20$ ).

	$\tau_{\text{intra}}$ [ms]					$d / d_{\text{exp}}$ [ $\mu\text{m}$ ]				$p_{\text{intra}}$
	50	100	200	500	1000	3 / 0.1	7 / 2.0	11 / 4.6	15 / 8.3	
Simulated	50	100	200	500	1000	3 / 0.1	7 / 2.0	11 / 4.6	15 / 8.3	80%
Estimated	$45 \pm 4.6$	$89 \pm 8.0$	$180 \pm 11$	$490 \pm 4.7$	$1000 \pm 44$	$1.8 \pm 0.75$	$2.5 \pm 0.13$	$4.6 \pm 0.68$	$7.8 \pm 0.68$	$66 \pm 15\%$