

# Spectrally Selective Hard Pulses (SSHP) for Positive Contrast MRI: Theory and Validation

T. Gupta<sup>1</sup>, S. Shah<sup>2</sup>, S. Virmani<sup>3</sup>, R. Omary<sup>4,5</sup>, and A. Larson<sup>1,5</sup>

<sup>1</sup>Radiology and Biomedical Engineering, Northwestern University, Chicago, IL, United States, <sup>2</sup>Siemens Medical Solutions, Chicago, IL, United States, <sup>3</sup>Radiology and Biomedical Engineering, Northwestern University, <sup>4</sup>Radiology, Northwestern University, Chicago, IL, United States, <sup>5</sup>Robert H. Lurie Comprehensive Cancer Center, Northwestern University, Chicago, IL, United States

**Introduction:** Several approaches have been proposed for signal-enhanced visualization of paramagnetic markers and contrast agents used for catheter tracking, atherosclerotic plaque imaging and stem cell research [1-3]. Recently, Patil *et al* proposed using spatially selective binomial pulses to generate positive contrast from paramagnetic markers [4]. In this work we investigated the use of spectrally selective hard pulses (SSHP) for positive contrast MRI of paramagnetic markers. Our primary focus was to study the effects of specific SSHP-parameters upon marker conspicuity, background suppression and fat suppression.

**Theory: (I)** Susceptibility-shifted paramagnetic markers exposed to the static magnetic field ( $B_0$ ) generate a range of dipolar frequency shifts ( $\Delta\omega$ ) within the surrounding tissues. Conspicuity of a paramagnetic marker is directly proportional to the total enhanced surface area contributed by all excited off-resonance (off-res) iso-frequency surfaces surrounding the marker. The enhanced surface area ( $S_{enh}$ ) around a spherical paramagnetic marker decreases rapidly for iso-frequency surfaces with higher values of  $\Delta\omega$ :

$$S_{enh} = \frac{4\gamma\pi a^3(\Delta\gamma) B_0^*(0.7698)}{9(\Delta\omega^2)} \quad [\text{Eqn 1}]$$

$[a = \text{radius of the paramagnetic sphere, } \gamma = \text{gyromagnetic ratio}]$

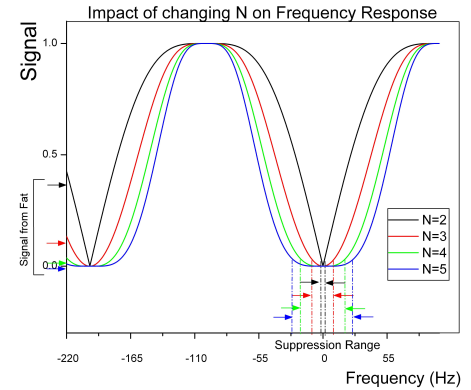
**(II)** A SSHP sequence contains a series of N short rectangular RF sub-pulses with flip angles  $\Phi_k(k=1 \text{ to } N)$  and  $\tau$  RF spacing (sec) between two adjacent sub-pulses. The applied RF phase alternates between 0 and  $\pi$  radians across the sub-pulse series. The SSHP sequence selectively excites specific off-res targets ( $\Delta\omega$ ) around the paramagnetic marker. These targets can be selected by varying the RF spacing ( $\tau$ ):  $\tau = (1/4*\pi*\Delta\omega)$ . In addition to exciting the targeted iso-frequency surface, the SSHP pulse also excites frequencies at  $(\Delta\omega+\Delta\omega*i)$  ( $i=2,4,6,8..$ ) while suppressing the on-resonant (on-res) frequencies and all off-res frequencies at  $(\Delta\omega*i)$ .

**Materials and Methods:** A standard GRE sequence was modified to replace the conventional excitation scheme with composite SSHP pulses using 24 different combinations of N (# of sub-pulses) and  $\tau$  (RF spacings). For each N=2, 3, 4 and 5,  $\tau$  was adjusted to target a range of off-res frequencies including 100, 200, 300, 400, 800, 1600 Hz, respectively. The sub-pulse flip angles were individually weighted by binomial pulse ratios of order (N-1) to generate an effective flip angle of 90°. SSHP images of an agar phantom with an imbedded spherical paramagnetic  $Fe_2O_3$  particle were acquired using a Siemens Sonata 1.5T clinical scanner. Sequence parameters: FOV 250x172 mm, Matrix 250x172, Slice/Slab = 16, Slice thickness = 2mm, TR/TE = 80/(2.82ms + (N-1)\* $\tau$ ), BW = 260Hz/Px. MR images and corresponding MATLAB simulations were evaluated for conspicuity, background suppression and fat suppression.

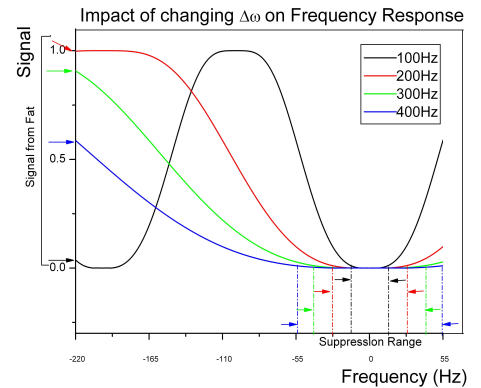
**Results:** Each of the 24 N- $\tau$  combination SSHP pulses generated a characteristic dipolar signal enhancement pattern around the paramagnetic  $Fe_2O_3$  particle (Fig.3). SSHP pulses with higher N values resulted in a broader suppression range (Fig.1), as a result on-res protons and  $B_0$  inhomogeneities were more effectively and homogeneously suppressed at higher N values (Fig.3). Suppression range also increased with increasing  $\Delta\omega$  (Fig.2). Figure 3 shows rapidly declining conspicuity with increasing N and increasing off-res target  $\Delta\omega$ . This is due to (a) increasing suppression range at higher N and  $\Delta\omega$  values and (b) decreasing  $S_{enh}$  with increasing  $\Delta\omega$  [Eqn 1]. According to MATLAB simulations, fat suppression increased with increasing N due to overall increased background suppression (Fig.1). Furthermore, for  $\Delta\omega=110$ Hz, frequencies around 220Hz ( $\Delta\omega*[i=2]$ ) should be completely suppressed in addition to on-res protons as described above. Amongst our selected  $\Delta\omega$  targets, 100 Hz was closest to 110 Hz, consequently fat suppression was maximum at  $\Delta\omega =100$ Hz (Fig.2). According to graphical simulations, fat signal was slightly lower for  $\Delta\omega =1600$  Hz. However,  $S_{enh}$  (being inversely proportional  $\Delta\omega^2$ ) was minimal at 1600Hz resulting in very low signal contribution from these protons and consequently, poor image quality due to low SNR within these images (Fig.3).

**Discussion:** Incorporation of SSHP pulses in standard GRE sequences offers a promising method for positive contrast MRI. N,  $\tau$  and  $\Delta\omega$  may be individually optimized for applications requiring high fat suppression (e.g. abdominal imaging), homogenous background suppression (e.g. high-field imaging) or greater conspicuity (e.g. iron-labeled cell tracking). Further developments are needed to demonstrate the utility of this technique for *in-vivo* clinical applications. **Future directions:** (1) TSE based application of binomial hard pulses, (2) addition of spectrally selective FATSAT pulse to further saturate fat signal [4], (3) application of non-exact binomial flip angles for increased background suppression at lower N values [5].

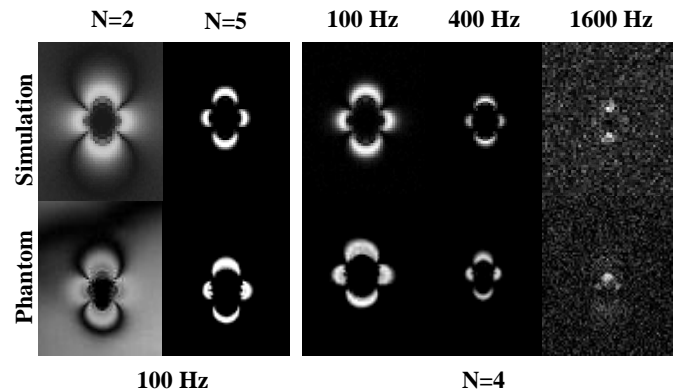
**References:** [1] Seppenwoolde *et al*, Magn Reson Med 50:784-790 (2003) [2] Stuber *et al*, Proc Intl. Soc. Mag. Reson. Med. 13(2006) [3] Mani *et al*, Magn Reson Med 56:1096-1106 (2006) [4] Patil *et al*, Proc Intl. Soc. Mag. Reson. Med. 14(2007) [5] Schick, MAGMA 1: 158-168 (1993)



**Fig. 1** Frequency response of SSHP at off-res target of 100Hz for series of N = 2, 3, 4, 5 pulses. The suppression range broadens (dashed lines) and the fat signal decreases (colored arrows marking the signal axis at frequency of -220 Hz) with increasing N value.



**Fig. 2** Frequency response of SSHP with N=4 for  $\Delta\omega =100, 200, 300, 400$  Hz. The suppression range broadens (vertical dashed lines) with increasing  $\Delta\omega$ . The fat signal (colored arrows marking the signal axis at frequency of -220 Hz) is lowest for  $\Delta\omega = 100$ Hz.



**Fig. 3** SSHP simulation (top row) and agar phantom (bottom row) images of  $Fe_2O_3$  particle. For a constant off-res target=100Hz, background suppression increases and conspicuity decreases as N increases from 2-5 (from the left: col. 1 and 2). For a constant N=4, conspicuity decreases rapidly with increasing off-res target value (col.3, 4 and 5). Low  $S_{enh}$  at 1600 Hz iso-frequency surface leads to low signal contribution (col. 5).