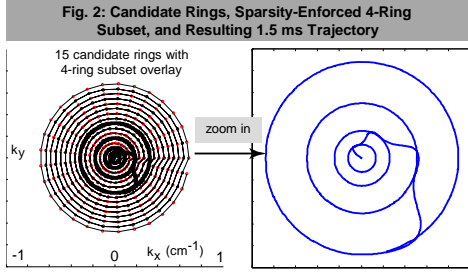
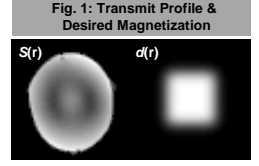


# Sparsity-Enforced Joint Spiral Trajectory & RF Excitation Pulse Design

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**INTRODUCTION.** Algorithms have recently been proposed that jointly optimize excitation  $k$ -space trajectories & RF excitation pulses [1-4]. Each relies on a small-tip-angle assumption and has been shown to produce higher-quality excitations or shorter excitation pulses than conventional 3D spoke-based [1,2], 2D echo-planar (EP) [3], and 2D spiral trajectories [4]. Here we propose a sparsity-enforcement algorithm that jointly determines sparse, quickly-traversable  $k$ -space trajectories and excitation pulses. The method is applicable to single- as well as multi-channel transmission systems. An  $L_1$  penalty is used when searching over a large number of possible trajectory segments to reveal a small subset of these segments (along with an RF pulse) that *alone* form a high-fidelity version of a user-specified target excitation. These segments may be



of any shape or size, but here, we provide the algorithm with *candidate rings* in  $k$ -space, thus focusing on the optimization of spirals. Simulations are conducted using data from a single-channel 7T system: for a Gaussian-shaped target and inhomogeneous transmit profile, we show that sparsity-enforced spirals lead to improvements in both excitation quality & pulse duration relative to conventional radially-undersampled (“accelerated”) spirals.

**METHODS. Transmission profile & desired target magnetization pattern.** Fig. 1 depicts the transmit profile,  $S(\mathbf{r})$ , of a birdcage coil in a water phantom at 7T, along with the desired magnetization pattern,  $d(\mathbf{r})$ . Forming  $d(\mathbf{r})$  in the presence of the inhomogeneous  $S(\mathbf{r})$  is non-trivial.

**Conventional spirals.** Consider a spiral whose radii are spaced at the Nyquist-limit (as defined by the desired FOV). Defining an unaccelerated Nyquist-sampled spiral an “ $R=1$ ” spiral, the conventional way to sparsify (accelerate) the trajectory is to radially undersample its rings by a factor of  $R$ . This process does indeed

accelerate the trajectory and reduce pulse duration, but does not explicitly take the desired magnetization pattern into account.

**Sparsity-enforced spiral design.** We develop a spiral acceleration method that accounts for  $d(\mathbf{r})$  up-front and determines *target-specific* trajectory-pulse combinations.

We start by deriving the algorithm for a single-channel system. Imagine  $C$  contours (e.g., ellipses, rings) in  $k$ -space; for each contour  $c$ ,  $N_c$   $k$ -space locations along the contour are known:  $\mathbf{k}_{c,1}, \dots, \mathbf{k}_{c,N_c}$ . Overall,  $N_k = N_1 + \dots + N_C$  discrete  $k$ -space points are known. The goal is to find the *smallest contour subset* along with a corresponding pulse such that when the pulse is played along the path defined by the small set of contours, a high-fidelity version of  $d(\mathbf{r})$  is produced; the contour segments in the small subset may then be connected via gradients to yield a highly-sparse, fast  $k$ -space trajectory. We pose an optimization that seeks an energy weighting to place in  $k$ -space to form  $d(\mathbf{r})$ , but where a penalty is incurred whenever a contour experiences a nonzero energy deposition. This ensures only a small number of contours will be used. First, the linearized formalism of [5] is applied, allowing us to relate energy placed along each  $k$ -space contour to the resulting magnetization,  $m(\mathbf{r})$ :  $\mathbf{m} = \mathbf{S}\mathbf{C}\mathbf{g}_c + \dots + \mathbf{S}\mathbf{C}_c\mathbf{g}_c = \mathbf{A}_{\text{tot}}\mathbf{g}_{\text{tot}}$  (Eq. 1), where  $\mathbf{m} \in \mathbb{C}^{M_s}$  contains samples of  $m(\mathbf{r})$  at  $\mathbf{r}_1, \dots, \mathbf{r}_{M_s}$ ,  $\mathbf{S} \in \mathbb{C}^{M_s \times M_s}$  is a matrix of  $S(\mathbf{r})$  samples, &  $\mathbf{g}_c \in \mathbb{C}^{N_c}$  contains weights the transmit channel places at the  $N_c$   $k$ -space locations of contour  $c$ . Each  $C_c$  relates how energy placed along contour  $c$  affects the resulting excitation: formally,  $C_c(u,v) = j\gamma\Delta_t M_0 \exp(j\mathbf{r}_u \cdot \mathbf{k}_{c,v})$ ,  $\mathbf{C} \in \mathbb{C}^{M_s \times N_c}$ ,  $\mathbf{A}_{\text{tot}} = [\mathbf{S}\mathbf{C}_1 \dots \mathbf{S}\mathbf{C}_C]$ , &  $\mathbf{g}_{\text{tot}} = [\mathbf{g}_1 \dots \mathbf{g}_C]^T$ . The desired excitation,  $d(\mathbf{r})$ , is vectorized to  $\mathbf{d} \in \mathbb{C}^{M_s}$ . Solving  $\mathbf{d} = \mathbf{A}_{\text{tot}}\mathbf{g}_{\text{tot}}$  (Eq. 2) for  $\mathbf{g}_{\text{tot}}$  via the pseudoinversion of  $\mathbf{A}_{\text{tot}}$  results in a solution where energy is placed along *all* contours, i.e., all  $\|\mathbf{g}_c\|_2$ s are nonzero, failing to reveal a useful contour subset. Consider however solving Eq. 2 while penalizing each nonzero contour energy (each  $\|\mathbf{g}_c\|_2$ ), which prohibits the use of many contours while encouraging those that remain to form  $\mathbf{g}_s$  that approximately solve Eq. 2. This is accomplished by solving  $\min_{\mathbf{g}} \{ \|\mathbf{d} - \mathbf{A}_{\text{tot}}\mathbf{g}_{\text{tot}}\|_2^2 + \lambda \sum_c \|\mathbf{g}_c\|_2 \}$  (Eq. 3) for fixed  $\lambda$ , where  $\sum_c \|\mathbf{g}_c\|_2$  is the  $L_1$ -norm of the  $\|\mathbf{g}_c\|_2$  contour energies; such a norm encourages sparsity [6,7]. As  $\lambda:0 \rightarrow 1$ , increasing numbers of contours have their energies driven to zero, residual error increases, and smaller contour subsets are revealed. For  $N_k < 1000$ , finding a small contour subset takes only several minutes. After solving Eq. 3 for a large  $\lambda$ , an optimized  $Q$ -contour subset is formed, comprised of those  $Q$  contours whose  $\|\mathbf{g}_c\|_2$ s are largest. The  $Q$  segments are then connected via an in-house method, yielding  $\mathbf{k}(t)$  &  $\mathbf{G}(t)$ . Finally, the RF pulse is returned by truncating  $\mathbf{A}_{\text{tot}}$  &  $\mathbf{g}_{\text{tot}}$  and least-squares fitting to Eq. 2. For a  $P$ -channel parallel transmission system,  $\mathbf{A}_{\text{tot}}$  of Eq. 1 will contain  $PC$  rather than  $C$  submatrices, there will be  $PC$   $\mathbf{g}_{p,c}$  weight vectors, and Eq. 3’s regularization term becomes  $\sum_c \{ \|\mathbf{g}_{1,c}^T \dots \mathbf{g}_{p,c}^T\|_2 \}$ , i.e., only *overall* contour energy is penalized; whether it is one channel or all channels that make large contributions to a contour does not matter if the use of this contour greatly helps at forming the excitation.

Fig. 3: Sparsity-Enforced 4-Ring RF Pulse and Gradients

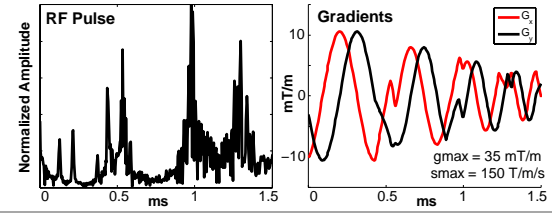
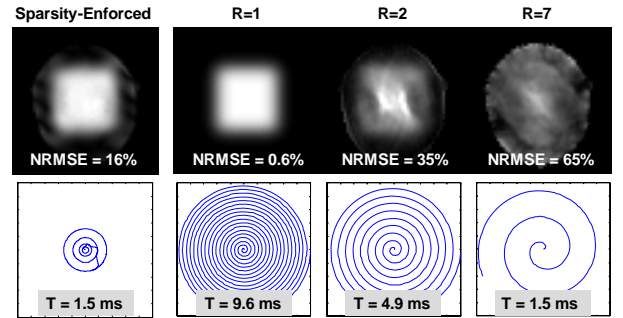


Fig. 4: Excitations due to Sparsity-Enforced & Accelerated Spirals



**RESULTS.** A sparsity-enforced spiral trajectory and RF pulse are designed by first defining the  $C = 16$  contours shown in the left panel of Fig. 2; there are 15 candidate rings, each comprised of 25  $k$ -space points, along with a single point at DC (thus  $N_k = 376$ ). These contours, along with  $S(\mathbf{r})$  and  $d(\mathbf{r})$ , are provided to the method,  $\lambda$  is set to 0.30, and Eq. 3 is then solved in under 2 minutes. Four rings are retained as the sparsity-enforced subset and appear as overlays on the left panel of Fig. 2. The right panel of Fig. 2 shows how the 4-ring subset is connected into a 1.5-ms trajectory, and Fig. 3 depicts the corresponding gradients and retuned RF pulse. Fig. 4 compares excitations due to the sparsity-enforced spiral and  $R$ -accelerated conventional spirals. The 1.5-ms sparsity-enforced spiral significantly outperforms the 4.9-ms,  $R = 2$  spiral. The  $R = 1$  spiral does indeed produce a near-perfect excitation, but is 6.4x longer than the optimized pulse. Note how the sparsity-enforced spiral traverses only a small segment of  $k$ -space yet is capable of forming a high-fidelity version of the box, in spite of the presence of the inhomogeneous  $S(\mathbf{r})$ . When  $R = 7$ , the conventionally-accelerated spiral has a duration on-par with the optimized spiral, but has 4x larger error. Thus for fixed excitation quality, the sparsity-enforced spiral yields shorter pulses, and for fixed pulse duration, it yields lower-NRMSE excitations.

**CONCLUSION.** A spiral trajectory acceleration algorithm based on sparsity-enforcement concepts has been shown to rapidly calculate fast, high-quality trajectories & corresponding RF pulses. Sparsity-enforced spiral trajectories significantly outperformed conventional spirals in simulated trials in a single-channel, 7T, non-uniform transmit profile setting.

**REFERENCES.** [1] Setsompop, Zelinski et al. ISMRM ’07, p 356. [2] Zelinski et al. ISMRM ’07, p 1691. [3] Yip et al. MRM ’07:58(3):598-604. [4] Levin et al. ISMRM ’06, p 3012. [5] Grissom et al. MRM. 2006;56(3):620-629. [6] Chen et al. SIAM Rev ’01;43(1):129–159. [7] Malioutov et al. IEEE Trans Sig Proc ’05;53(8).