

# Adiabatic $B_1$ Mapping for RF Current Density Imaging

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**Introduction:** RF current density imaging [1] can be used to predict RF ablation treatment patterns and to monitor the effectiveness of pacemaker insulators at RF frequencies. When the current is near the Larmor frequency, its associated magnetic field acts as a  $B_1$  field, exciting magnetization. Due to the large variations and intensity of the local field and main field inhomogeneities near wires, typical  $B_1$  measurement methods are impractical. We propose an adiabatic  $B_1$  mapping pulse that covers a large dynamic range and compensates for  $B_0$  variations.

**Theory & Methods:** Adiabatic excitation pulses are typically used because they create the same final flip angle for a wide range of  $B_1$  magnitudes. However, if the end frequency of the pulse is not increased all the way to on-resonance, a  $B_1$  magnitude dependent excitation results that can be used to map the  $B_1$  field variation. The angle of the excitation and therefore the transverse magnetization is dependent on the  $B_1$  strength, as shown in Figures 1 and 2. A large range of  $B_1$  values can be distinguished from a single excitation.

A  $90^\circ$  adiabatic excitation is used as a reference to remove the dependence on the strength of the initial magnetization. The  $90^\circ$  excitation is also used to map the phase of  $B_1$ . Accurate phase mapping is particularly important for current density imaging.

Off-resonance acts as an additional frequency offset in the adiabatic pulse, changing the excitation frequency. For two different excitations with positive and negative excitation frequency offsets, the off-resonance will have the opposite effect on each and can be compensated.

The adiabatic pulse is a modified sech/tanh pulse, adapted using the method of Ugurbil et al [3] to cover a 16X dynamic range of  $B_1$  and off-resonance of  $\pm 400$  Hz. A typical ablation current of 1 A<sub>RMS</sub> will create fields ranging from 0.14 to 1.4 G at distances between 1 mm and 1 cm from the wire. The pulse was optimized for this range, covering  $B_1$  magnitudes from 0.1 to 1.6 G. Bloch simulations were performed for off-resonances from -400 to 400 Hz,  $T_1$  of 500 ms, and  $T_2$  of 50 ms with additive Gaussian noise. The double angle method for  $B_1$  mapping [4] was also simulated for comparison with a 1.6 G pulse corresponding to a  $90^\circ$  rotation. A double angle sequence to fully cover the same range as the adiabatic pulse requires a series of 5 pulses from  $\alpha$  to  $2^4\alpha$  [5], two more than the adiabatic sequence.

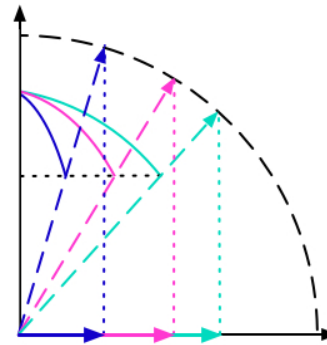
To test the feasibility of using an adiabatic pulse to measure  $B_1$  amplitude, experiments were run on a 1.5T GE Signa scanner with a phantom consisting of a wire with one inch exposed metal inserted into a tofu sample. Instead of using the  $90^\circ$  adiabatic excitation, the initial magnetization was assumed to be constant throughout the sample.

**Results:** The root-mean-squared-error of the  $B_1$  magnitude simulation results is shown in Figure 3. The double-angle pulse and the adiabatic pulse have similar accuracy for higher  $B_1$  values, but the adiabatic pulse covers the entire dynamic range with similar accuracy. Figure 4 shows the simulated accuracy of the  $B_1$  phase calculations. The performance of the adiabatic pulse suffers at low  $B_1$  magnitudes, but it performs better than the double angle method for the rest of the range, particularly in areas with large off-resonance effects.

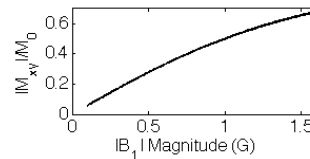
Figure 5 shows the reconstructed  $B_1$  magnitude for a projection along the length of the wire in the phantom. There is signal dropout near the wire due to cancellation of the fields on opposite sides of the wire, as the projection axis is not perfectly aligned with the wire. The field drops off rapidly away from the wire, as expected.

**Discussion & Conclusions:** RF safety requires the ability to measure current levels and leakage from implanted wires. RF ablation treatment patterns can be predicted with the knowledge of current distributions. Current-carrying wires create large variations in  $B_1$  magnitude near the wire while the presence of the wire creates inhomogeneity effects. The adiabatic  $B_1$  mapping pulse can cover a 16X dynamic range in 60% of the time of a double angle sequence and has better performance in measuring the  $B_1$  phase in the presence of off-resonance.

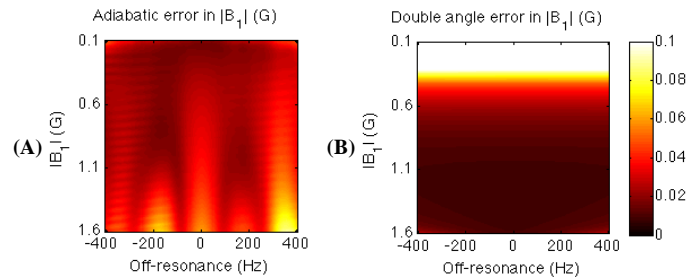
**References:** [1] G. Scott et al, Magn. Reson. Med. 28:186, 1992; [2] S. Conolly et al, J. Magn. Reson. 83:549, 1989; [3] K. Ugurbil et al, J. Magn. Reson. 80:448, 1988; [4] R. Stollberger et al, Proc 7th SMRM, p 106, 1988; [5] A. Kerr et al, Proc 15th ISMRM, p. 352, 2007



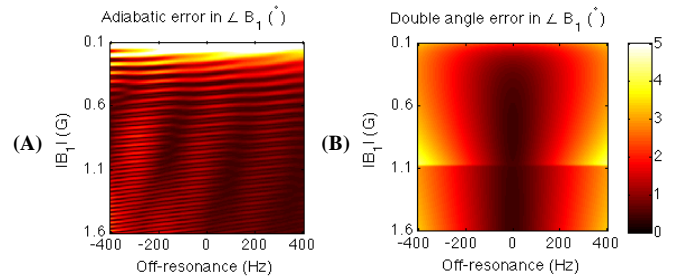
**Figure 1:** Sweep diagram [2] for adiabatic field mapping with various  $B_1$  magnitudes. Each color represents a different  $B_1$  magnitude. The solid arcs are the  $B_1$  vector paths. The black dotted line represents the frequency where the excitation ends. The colored dashed lines are the magnetization vectors corresponding to the excitation angles. The solid colored arrows are the final transverse magnetization.



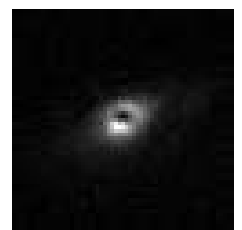
**Figure 2:** Transverse magnetization for adiabatic pulse with varying  $B_1$  magnitude. The magnetization vs. magnitude is approximately linear over the 0.1 to 1.6 G range.



**Figure 3:** Simulated error in  $B_1$  magnitude calculation (in Gauss) for the adiabatic method (A) and the double angle method (B). The adiabatic method covers a larger dynamic range without much loss in accuracy.



**Figure 4:** Simulated error in  $B_1$  angle calculation (in degrees) for the adiabatic method (A) and the double angle method (B). The adiabatic method is more accurate except at low  $B_1$  magnitudes.



**Figure 5:** Reconstructed  $B_1$  magnitude projection along length of wire in tofu phantom.