

Generalized Parameter Relations for the Shinnar-Le Roux Pulse Design Algorithm

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Introduction: The magnetization ripple amplitudes δ from a pulse designed by the Shinnar-Le Roux (SLR) algorithm are a non-linear function of the SLR A and B polynomial ripples δ . Pauly *et al* derived these relations for five types of pulses shown below [1]. However, these expressions do not cover the entire range of flip angles, e.g. neither the small-tip-angle or 90° relations for excitation pulses are exactly right for a flip angle of say, 70° . In deriving these relations, Pauly *et al* used a geometrical argument for excitation pulses, and a more general approach was used for other pulse types. Here, the general approach is applied to all the pulses, and gives the required general parameter relations.

Method: If the initial magnetisation is in the z -direction only, $M_{xy}^+ = 2A * BM_z^- (1)$, and $M_z^+ = (1 - BB^*)M_z^- (2)$, where A and B are the SLR polynomials of the rf pulse. If initial magnetisation is in the y -direction, after a spin-echo pulse surrounded by crusher gradients: $|M_{xy}^+| = |B^2 M_y^-| (3)$. In the design process, given the desired pulse duration, slice width, and ripples in and outside the passband (δ_1 and δ_2), the optimum transition bandwidth W is obtained using an empirical formula (Eq. 20 in Ref. [1]). Parameters δ_1 , δ_2 , and W are then passed to the Parks-McClellan algorithm which returns an equi-ripple polynomial, which is scaled to give the B polynomial. This also determines the magnitude of the A polynomial via: $|A|^2 + |B|^2 = 1 (4)$. Following scaling, the maximum (or minimum) and average A and B values are substituted into Eqs 1-4 and subtracted from each other to give the ripples δ , which can be graphed and inverted numerically to derive δ from a required δ .

Results: For all excitation pulses with flip angles ϕ not equal to 180° , $|\delta_1^e| = \left| \frac{2\sqrt{1 - (1 + \delta_1)^2 \sin^2(\phi/2)}(1 + \delta_1) \sin(\phi/2) \mp \sin \phi}{\sin \phi} \right| (5)$, and

$|\delta_2^e| = \left| \frac{2\sqrt{1 - \delta_2^2 \sin^2(\phi/2)}\delta_2 \sin(\phi/2)}{\sin \phi} \right| (6)$. The sign change in (5) accounts for change from $\phi < 180^\circ$ to $> 180^\circ$, and expressions have been normalised to the average in-slice magnetisation. Other expressions are:

Pulse type	Parameter of interest	Pauly relations		Exact relations	
		δ_1^e	δ_2^e	δ_1^e	δ_2^e
Small-tip angles	M_{xy}	δ_1	δ_2	Eq. 5	Eq. 6
90°	M_{xy}	$2\delta_1^2$	$\sqrt{2}\delta_2$	$\sqrt{2 - (1 + \delta_1)^2} (1 + \delta_1) - 1$	$\sqrt{2 - \delta_2^2} \delta_2$
Inversion	M_z	$8\delta_1$	$2\delta_2^2$	$8\delta_1 / (1 + \delta_1)^2$	$2\delta_2^2 / (1 + \delta_1)^2$
Crushed spin-echo	M_{xy}	$4\delta_1$	δ_2^2	$4\delta_1 / (1 + \delta_1)^2$	$\delta_2^2 / (1 + \delta_1)^2$
Suppression	M_z	$2\delta_1$	δ_2^2	$2\delta_1 + \delta_1^2$	δ_2^2

Fig 1 shows simulated in-slice profiles from a 70° pulse designed using Eq 5 (solid), compared with those designed using Pauly's small-tip (dotted) and 90° (dashed) relations. Fig 2 shows simulated in-slice profile for 110° pulse designed with Eq 5 (solid) vs Pauly's 90° relation (dashed). Design specifications: slice width 2 kHz, ripple amplitudes = 1% for both $\delta_{1,2}$, and pulse duration 4 ms. Horizontal lines show the required ripple, which is achieved by the generalized parameters. The generalized parameters also gave the required ripples outside the slice (not shown).

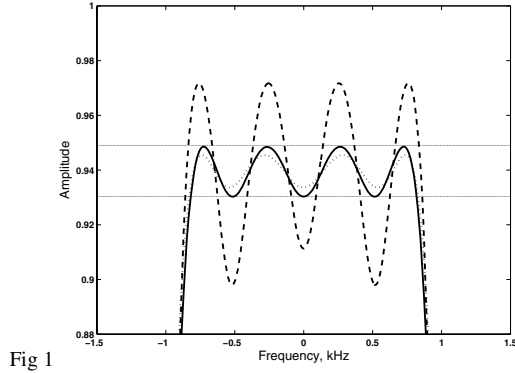


Fig 1

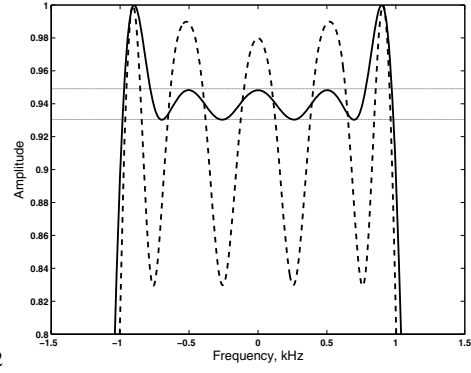


Fig 2

Discussion and Conclusion: Generalized parameter relations have been derived and show an improvement compared with the Pauly relations for flip angles not previously considered.

References: [1] Pauly J, Le Roux P, Nishimura D, Macovski A. IEEE Trans Med Imag (1991) 10:55-65.

Acknowledgments: EPSRC First Grant Scheme EP/C537491/1