

# Non-Fourier Reconstruction Can Prevent Image Dropouts

R. W. Cox<sup>1</sup>, L. A. McCall<sup>2</sup>

<sup>1</sup>National Institute of Mental Health, Bethesda, MD, United States, <sup>2</sup>Dept of Physics, Auburn University, Auburn, Alabama, United States

## Introduction:

Uncompensated through-slice magnetic field gradients cause extensive signal dropouts in full  $k$ -space rectangular-scan EPI due to the long interval before TE is reached. Techniques to fix this problem include  $z$ -shimming [1-2] and tailored RF pulses to alter the phase profile through the slice [3]. Herein, we propose another method, which in its simplest form does not need a modified EPI pulse sequence or any time penalty, just a new reconstruction algorithm. This method can also be used in conjunction with other techniques, since it only requires a magnetic field map and new software.

## The Signal-Magnetization Relationship and non-Fourier Reconstruction:

The basic problem is that the relationship between the measured signal and the magnetization is no longer precisely the Fourier transform. If we abstract out the  $x - k_x$  direction as being acquired too rapidly for dephasing to matter, then the signal at time  $t$  after excitation can be modeled as  $S(t) \propto \int_{-FOV/2}^{FOV/2} \int_{-h/2}^{h/2} M(y, z) e^{-ik_y(t)y} e^{-i[w(y)t+k_z(t)z]} e^{-R_2^*(y,z)t} dz dy$ , where  $M(y, z)$  is the post-RF transverse magnetization density (at  $t=0$ ), FOV is the field-of-view (in the  $y$ -direction),  $h$  is the slice thickness,  $k_y(t)$  is the phase-encoding readout,  $k_z(t)$  is the integral of any  $z$ -shimming gradients, and  $w(y)$  is  $\gamma$  times the spatially-dependent through-slice magnetic field gradient. Idealizing further that the magnetization profile in the  $z$ -direction is rectangular, the signal model becomes  $S(t) \propto \int_{-FOV/2}^{FOV/2} M(y) e^{-R_2^*(y)t} e^{-ik_y(t)y} \text{sinc}\left[\frac{h}{2}(w(y)t+k_z(t))\right] dy$ , where  $\text{sinc}(x) \equiv \sin(x)/x$ . If  $w(y)=0=k_z(t)$ , then this model is the usual Fourier relation between  $S(t)$  and  $M(y)$ . When  $k_z(t) \neq 0$ , at locations  $y$  where  $w(y)h \geq 2\pi/TE$ , the dephasing is such that  $\text{sinc}\left[\frac{h}{2}w(y)TE\right] \approx 0$  and virtually no signal is received from such values of  $y$  in the vicinity of  $t=TE$ ; normal Fourier reconstruction results in nearly total signal dropout in these locations. However, if the readout starts soon enough after excitation, the early part of the data *will* contain some signal from the dropout locations at nonzero values of  $k_y$  (while  $w(y)ht \leq 1.5\pi$ , say). The non-Fourier reconstruction model we

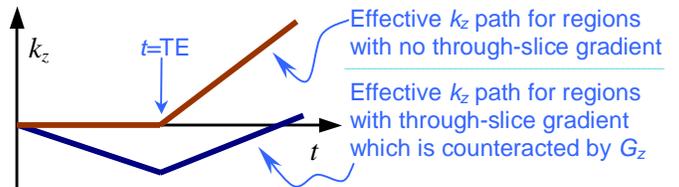
propose is thus to write the model for the image  $I(y)$  as  $S(t) \propto \int_{-FOV/2}^{FOV/2} I(y) e^{-ik_y(t)y} \text{sinc}\left[\frac{h}{2}(w(y)t+k_z(t))\right] dy$ , and then invert this linear integral equation numerically to solve for  $I(y)$ ; solving this equation requires having a  $B_z$  field map, so that  $w(y)$  is known. The discretization is  $S_m = \sum_n e^{-ik_y(t_m)y_n} \text{sinc}\left[\frac{h}{2}(w(y_n)t_m+k_z(t_m))\right] \cdot I_n$ , where  $n$  is the spatial index and  $m$  is the temporal index. This is a system of linear equations for the image vector  $[I_n]$  given the data vector  $[S_m]$ , and can be solved using the singular value decomposition (among many possible methods). As mentioned above, the speed of acquisition in the  $k_x$  direction means that normal FFT reconstruction can be used for the  $x - k_x$  inversion; therefore, although the matrix is dense, it can be solved quickly since it will involve (typically) only 64-128 unknown values. Once the subject- and orientation-specific gradient field map  $w(x, y) = \gamma \partial B_z(x, y, z) / \partial z|_{z=0}$  is available, the SVD for each slice and each  $x$  can be precomputed, and the overall reconstruction speed for a long imaging run of EPI data will be very fast.

## Preliminary Results and Discussion:

Simulated data were generated that correspond to a realistic EPI acquisition: 27 ms readout window starting at 6 ms after the center of the slice excitation RF; 64 data points in  $k_y$  with a 240 mm FOV; localized 100 Hz through-slice frequency change. The results show that knowing  $w(y)$  to within 10% provides an accurate image reconstruction.

Using a nonzero compensating  $k_z(t)$  for the second half of the acquisition provides a way to gain a second partial acquisition of data from “bad” regions. We keep  $G_z=0$  until  $t=TE$ ; for  $t>TE$ , a constant  $G_z$  is turned on to oppose  $w(y)$ .

Simulations show that this type of compensation, coupled with non-Fourier reconstruction, is less sensitive to errors in the estimated  $w(y)$  and to noise. Further work is planned to estimate the sensitivity to  $R_2^*$  changes and global field offsets, determine the PSF, choose the  $k_z(t)$  compensation path optimally, &c.



## References:

[1] Glover GH, *MRM* **42**: 290-299, 1999. [2] Song AW, *MRM* **46**: 407-411, 2001. [3] Pipe JG, *MRM* **33**: 24-33, 1995.